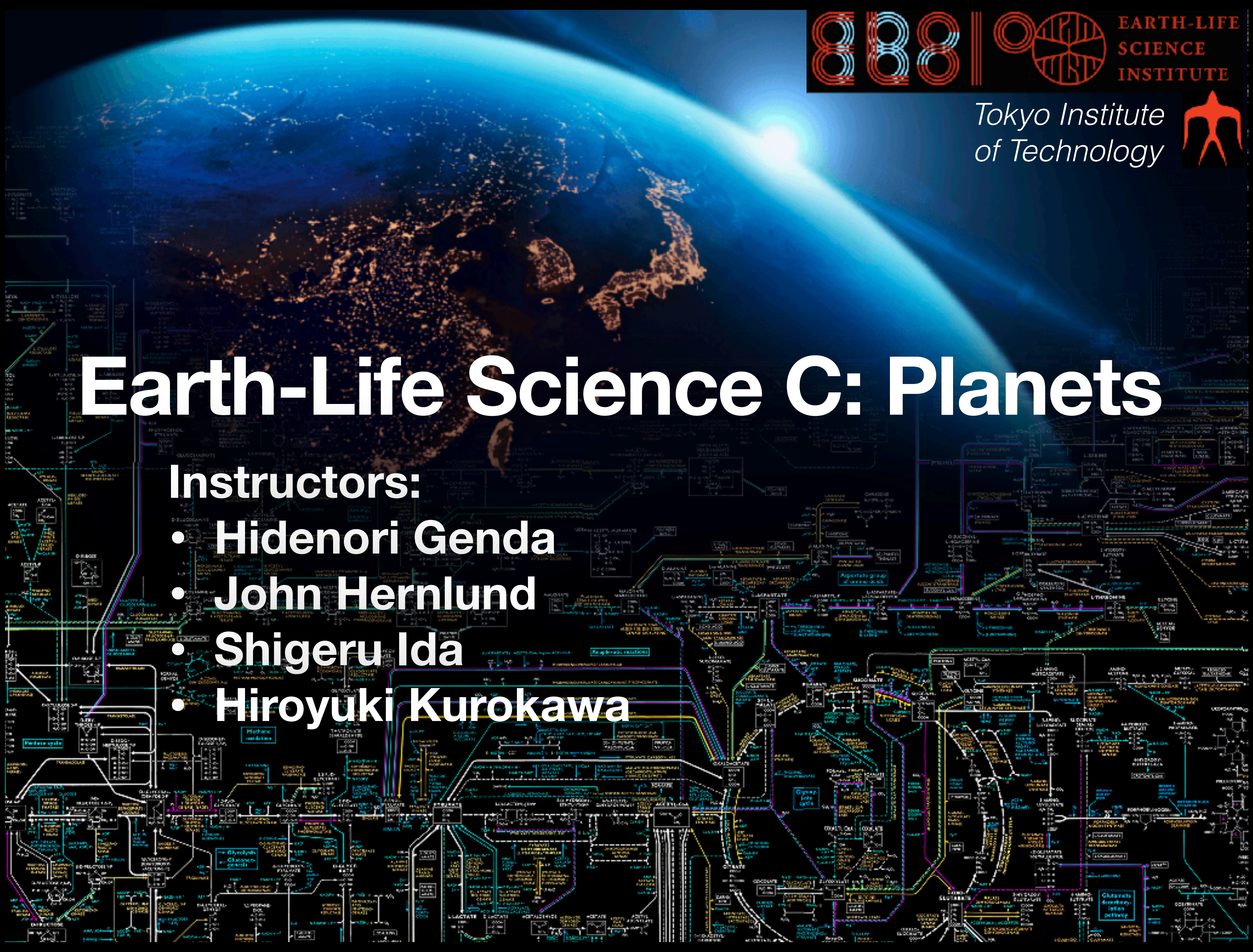


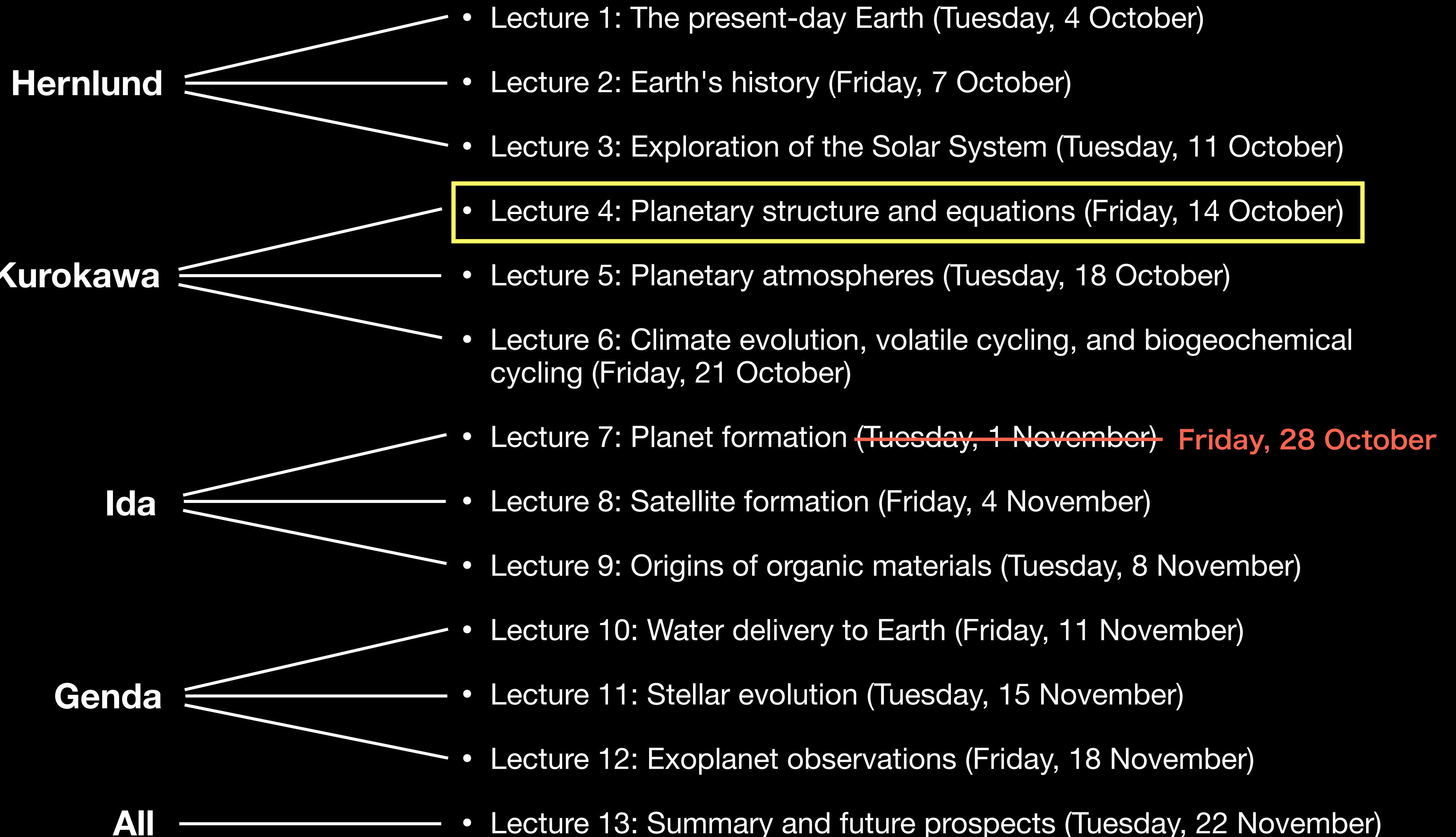


# Earth-Life Science C: Planets

Instructors:

- Hidenori Genda
- John Hernlund
- Shigeru Ida
- Hiroyuki Kurokawa



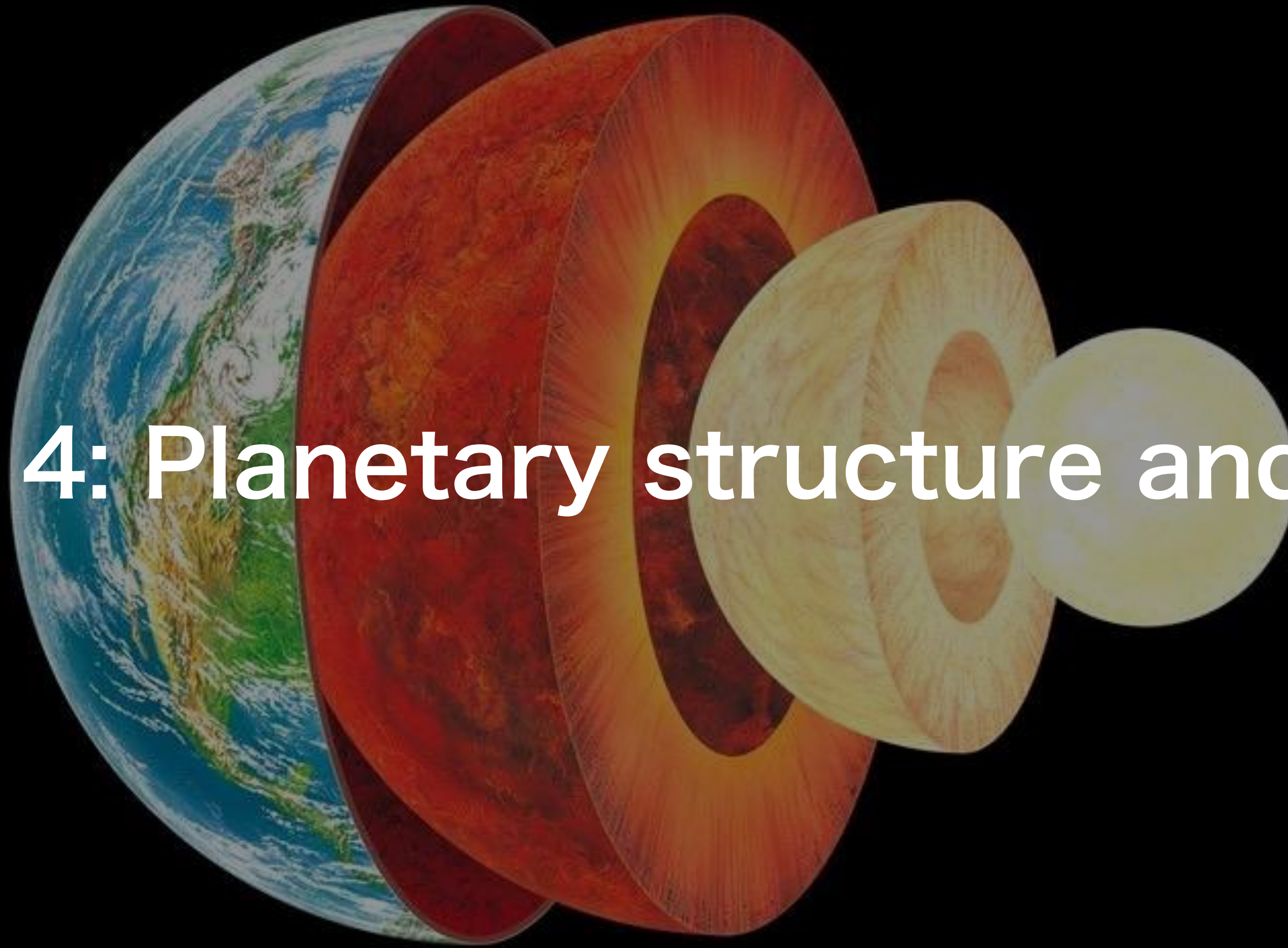


# About grade (score)

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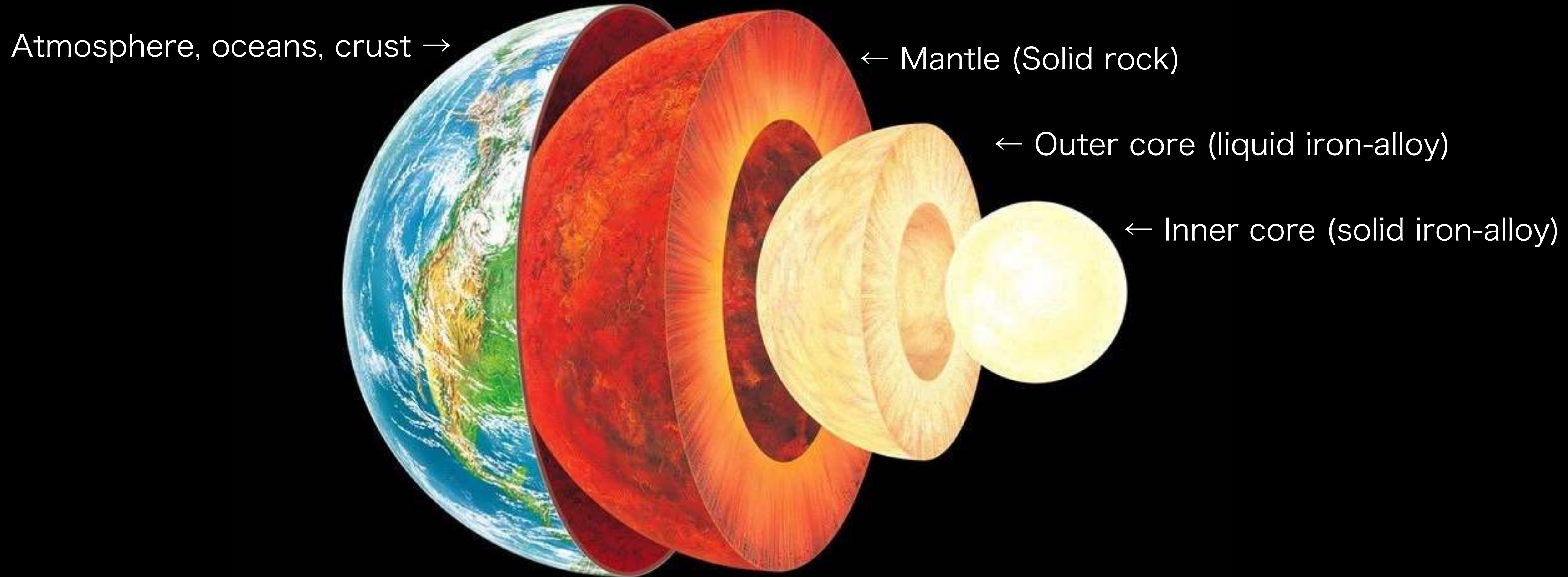
- Each lecturer will grade on a 25-point scale, and your grade will be the sum of the scores.
- The grade for my lectures will be evaluated based on weekly small reports. I will explain the assignments at the end of each lecture.

# Lecture 4: Planetary structure and equations



# Structure of Earth's interior

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- Rocky mantle + crust : 67.5 wt.%, Metallic core : 32.5 wt.%
- Atmosphere : Oceans : Solid Earth =  $8 \times 10^{-6} : 2 \times 10^{-4} : 1$

# Interiors of rocky planets

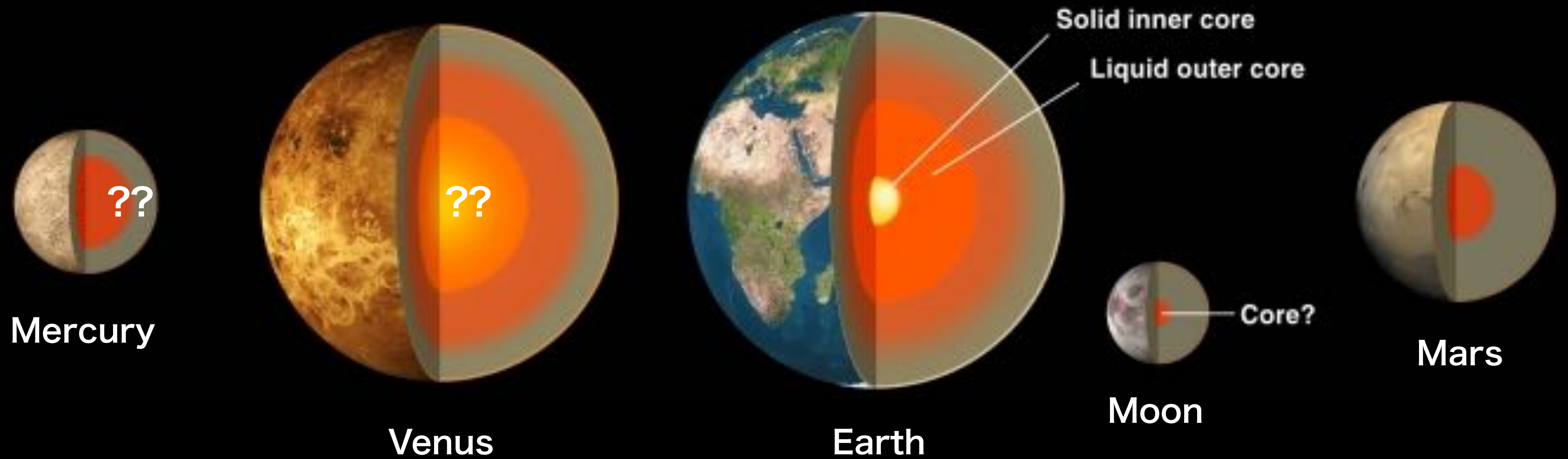


Image credit: NASA/LPI

# How do we know planetary interiors?

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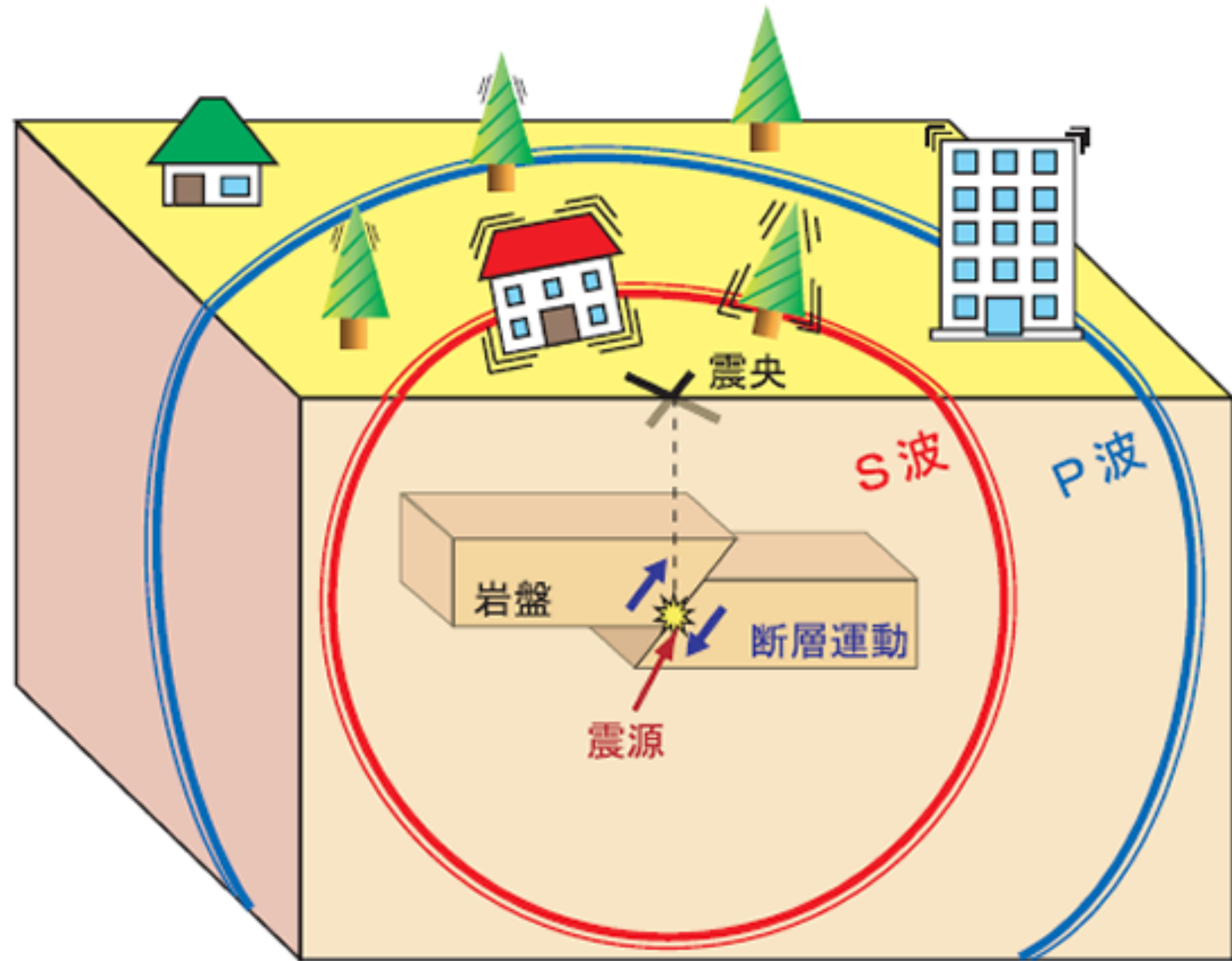
## Deep interior (> 10 km)

- Bulk density
- Magnetic field (dynamo, crustal remnants, induced)
- Moment of inertia (← precession, gravity field)
- Gravity field (higher order)
- Seismology
- Tidal deformation
- Analysis of materials originating from the deep interior (mantle xenolith)

## Shallow interior (< 10 km)

- Radar observations, gamma ray and neutron spectrometry, muography, etc.

# Seismology



## Speeds of primary and secondary waves $\Leftrightarrow$ Material properties

$$\text{P-wave : } v_p = \sqrt{\frac{K + 4\mu/3}{\rho}} \quad \text{--- (1)}$$

$$\text{S-wave : } v_p = \sqrt{\frac{\mu}{\rho}} \quad \text{--- (2)}$$

where  $\rho$ : density,  $\mu$ : shear modulus,  $K$ : bulk modulus

$$\mu \equiv \frac{F/A}{\Delta x/l} \quad \text{--- (3), } K \equiv -V \frac{\partial p}{\partial V} \quad \text{--- (4)}$$

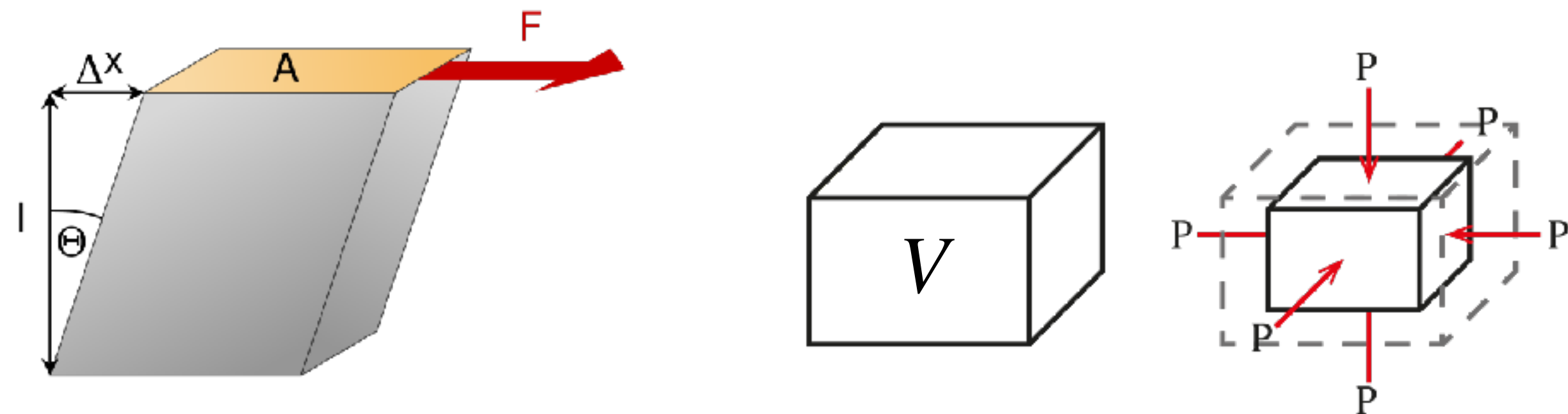
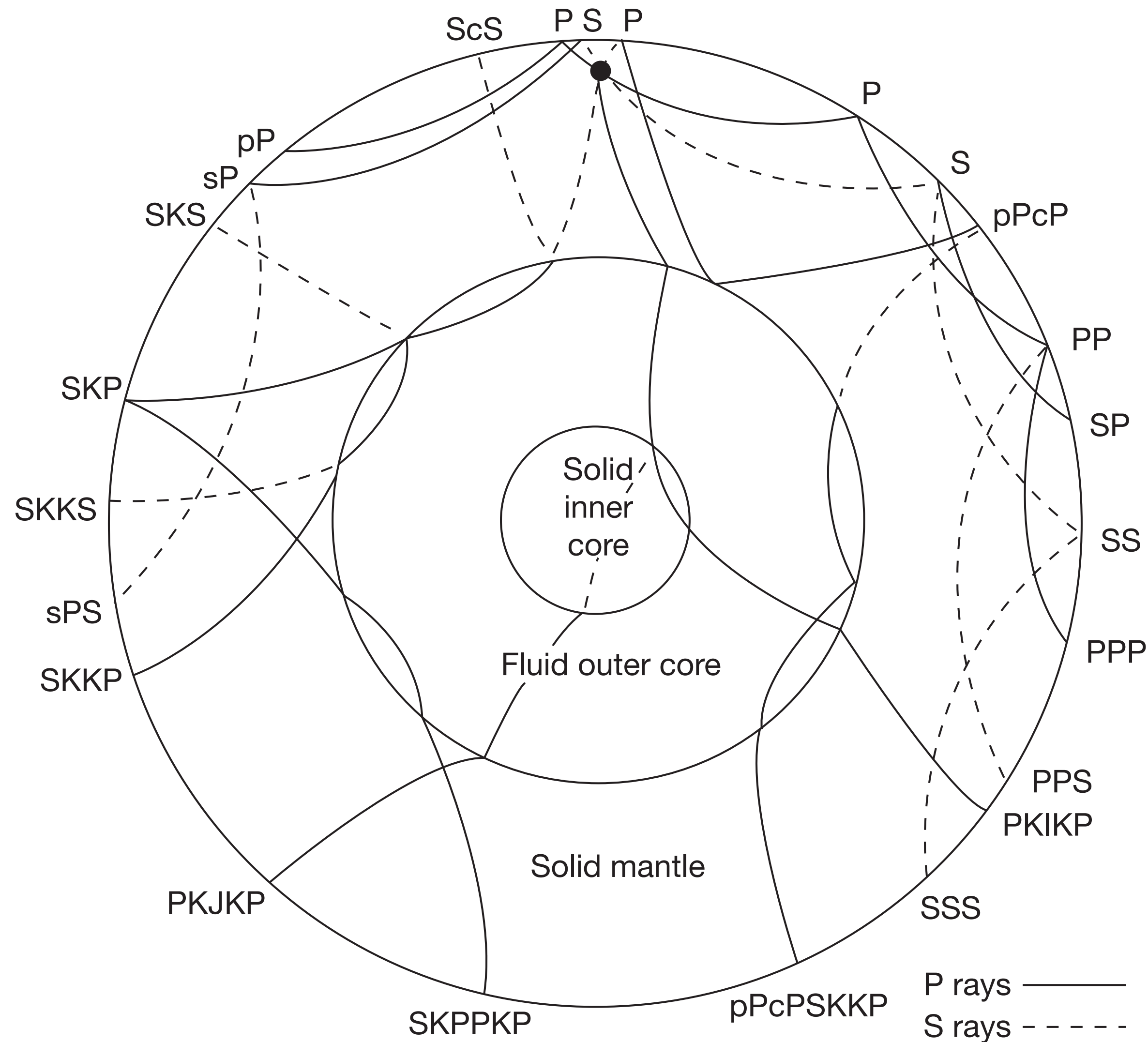


Image credit: 気象庁



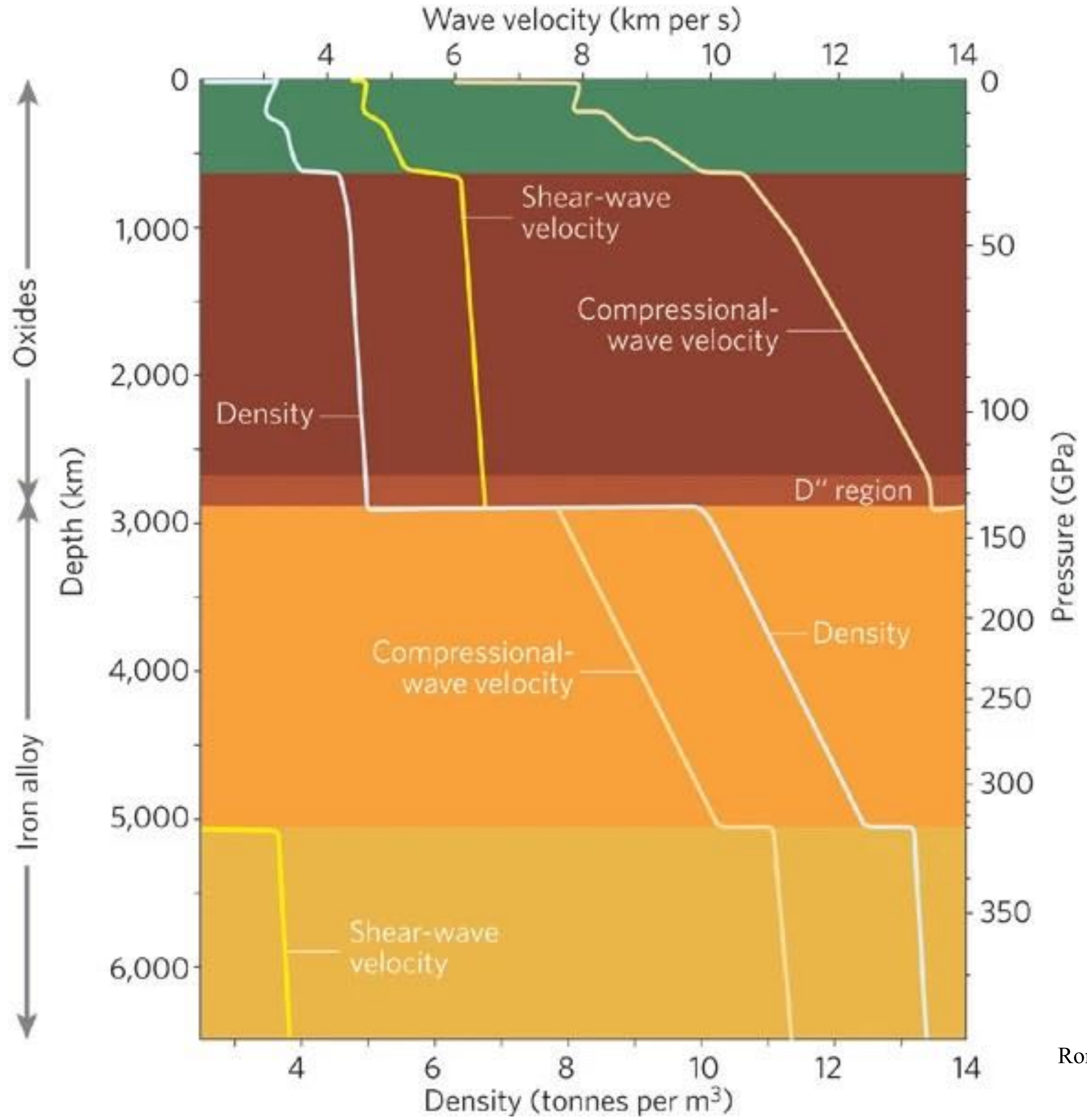
# Propagation of seismic waves in the interior



P: P-wave (solid), S: S-wave (dashed)  
 K: P-wave in the outer core, I: P-wave in the inner core,  
 J: S-wave in the inner core,  
 c: reflection at the core-mantle boundary.

**No S-wave in the outer core → liquid!**

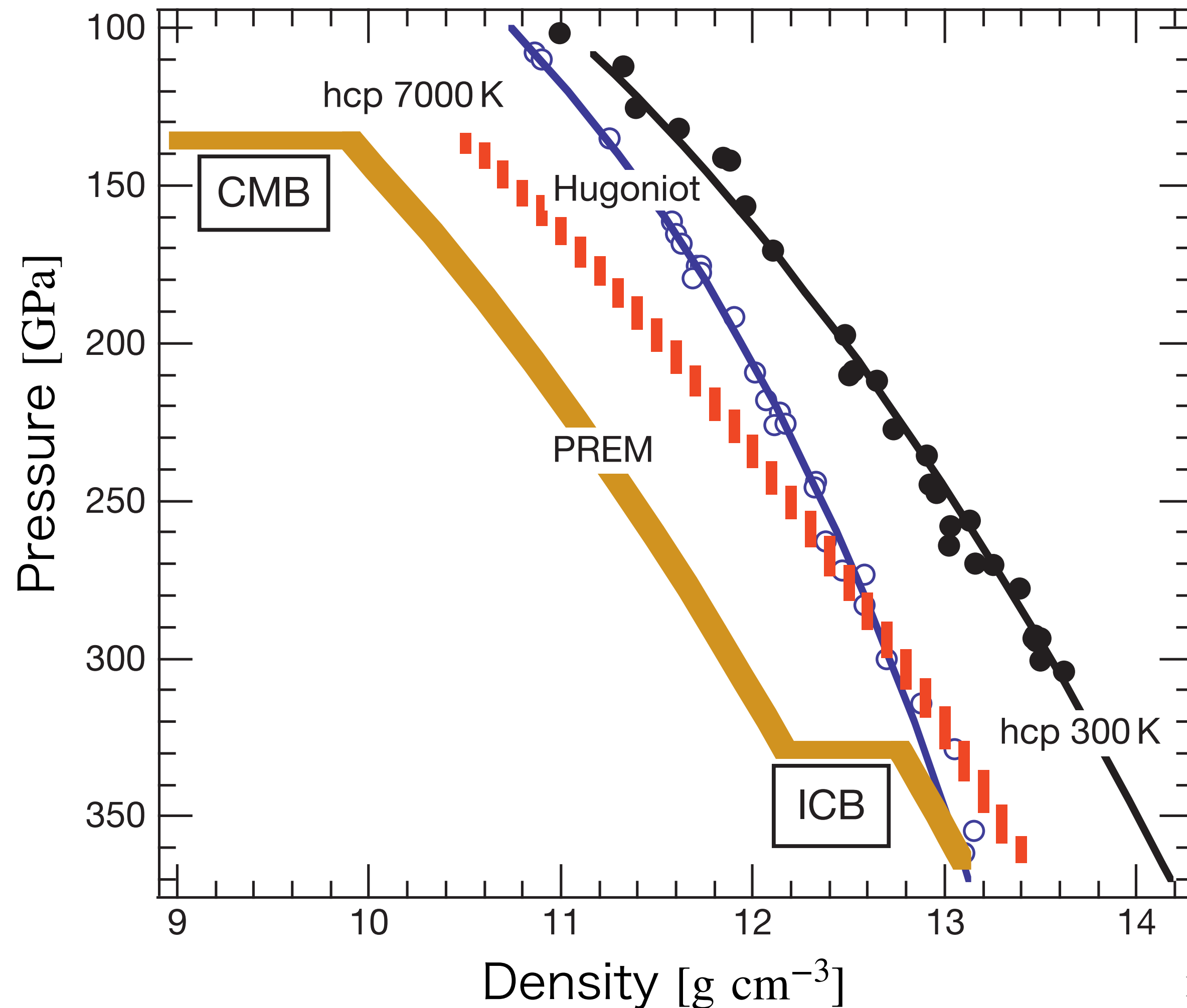
Dziewonski & Romanowicz (2015)  
 in *Treatise on Geophysics 2nd Edition*



Romanowicz (2008) *Nature*

# Light elements in the core

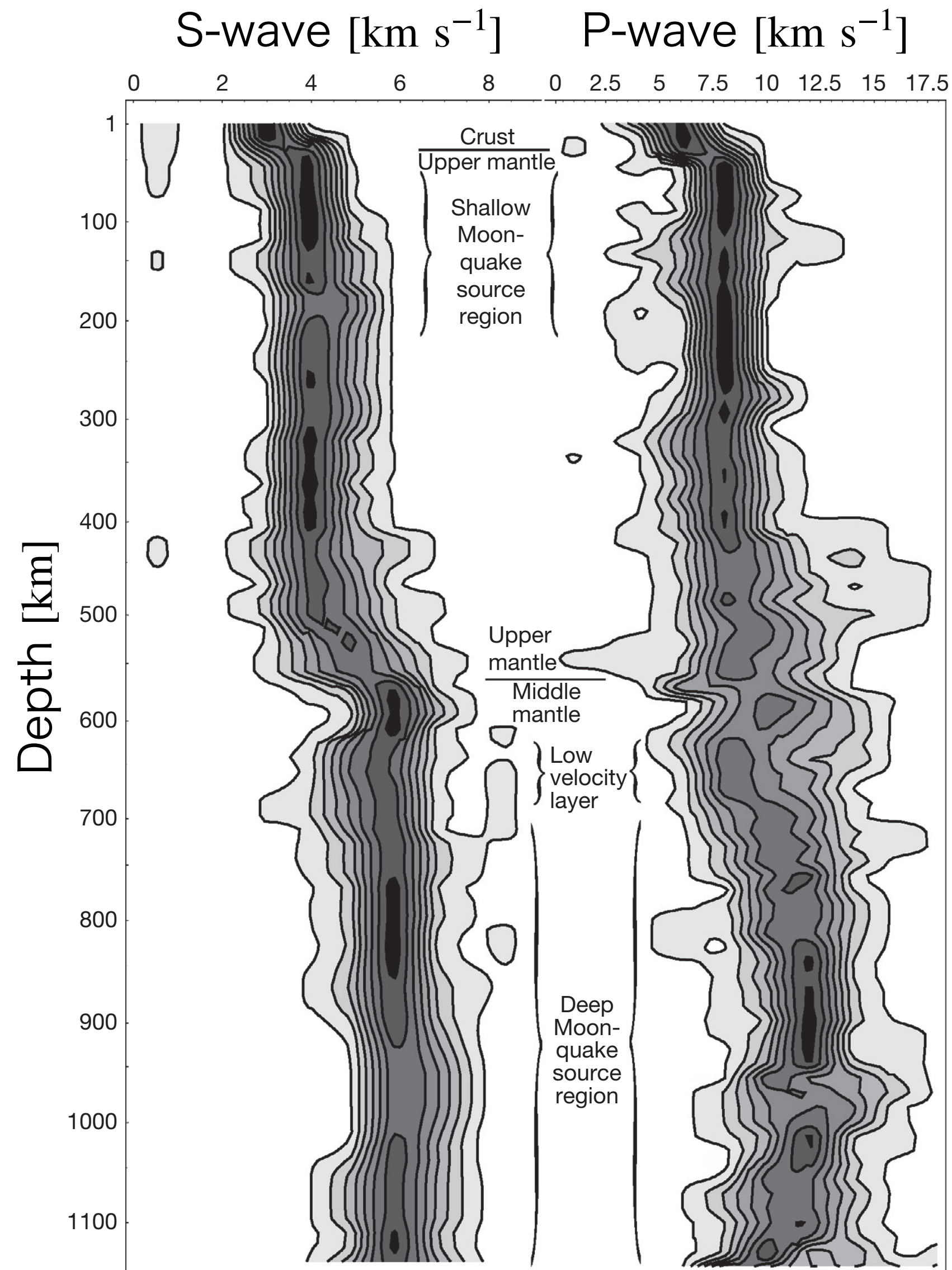
## Outer core density compared to that of pure Fe



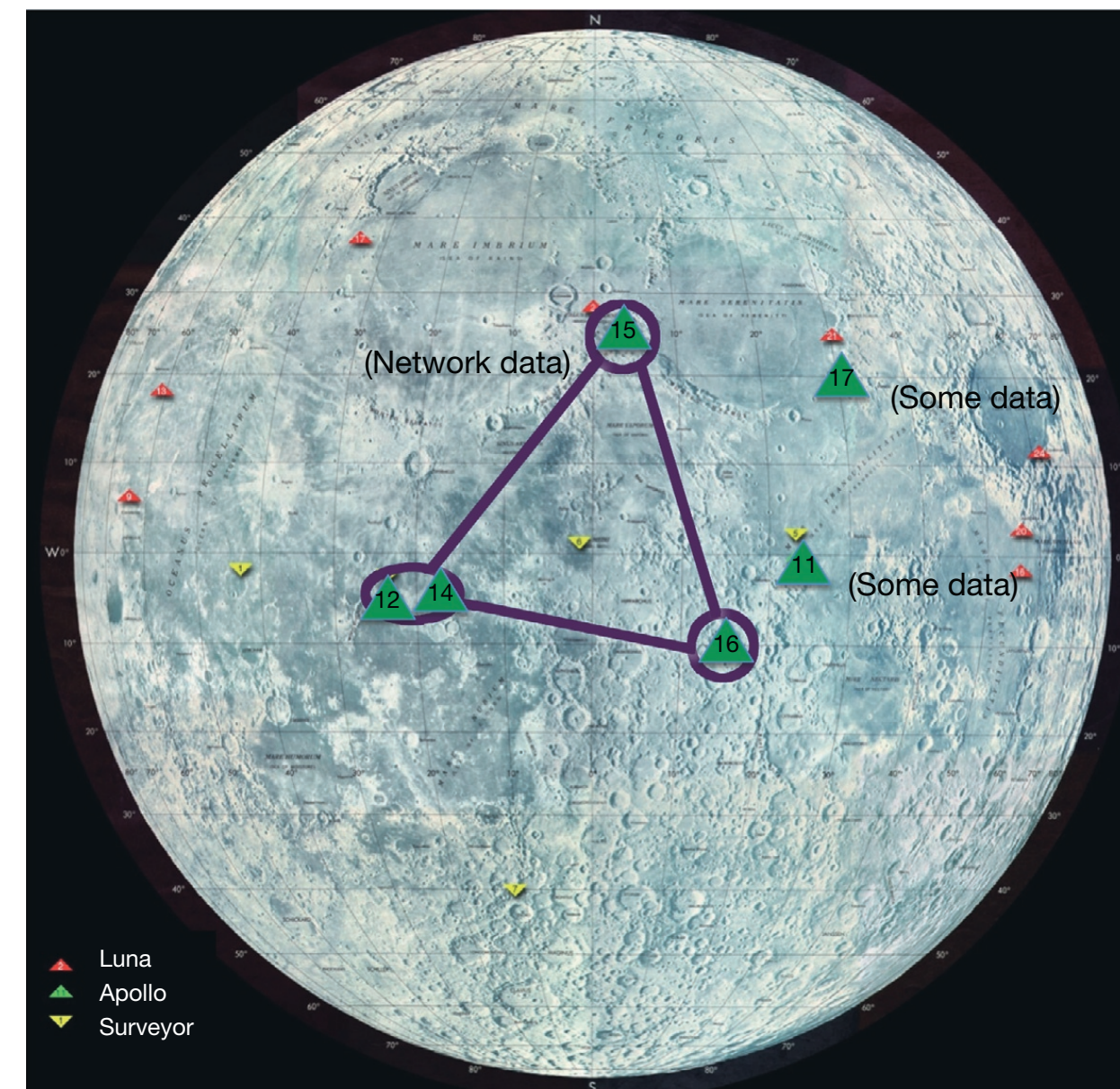
- The core is less dense than pure Fe by  $\sim 10\%$   
→ light element(s) – Si, S, O, H, and/or C?
- Total mass of the light elements in the core  
 $\sim M_{\oplus} \times 0.33 \times 0.1 \sim 2 \times 10^{22}$  kg  
 $\gg$  oceans ( $1.4 \times 10^{21}$  kg), atmosphere ( $5.1 \times 10^{18}$  kg)

Partitioning of light elements into the core has likely influenced determining Earth's surface environment

# Lunar seismology



- Seismograph network at Apollo 12, 14-16 landing sites
- Though data is limited compared to those on Earth, Lunar interior structure has been constrained by the seismology
  - Core size ~170–360 km (Nakamura et al. 1974)
  - Consistent with the estimates from the moment of inertia and induced magnetic fields

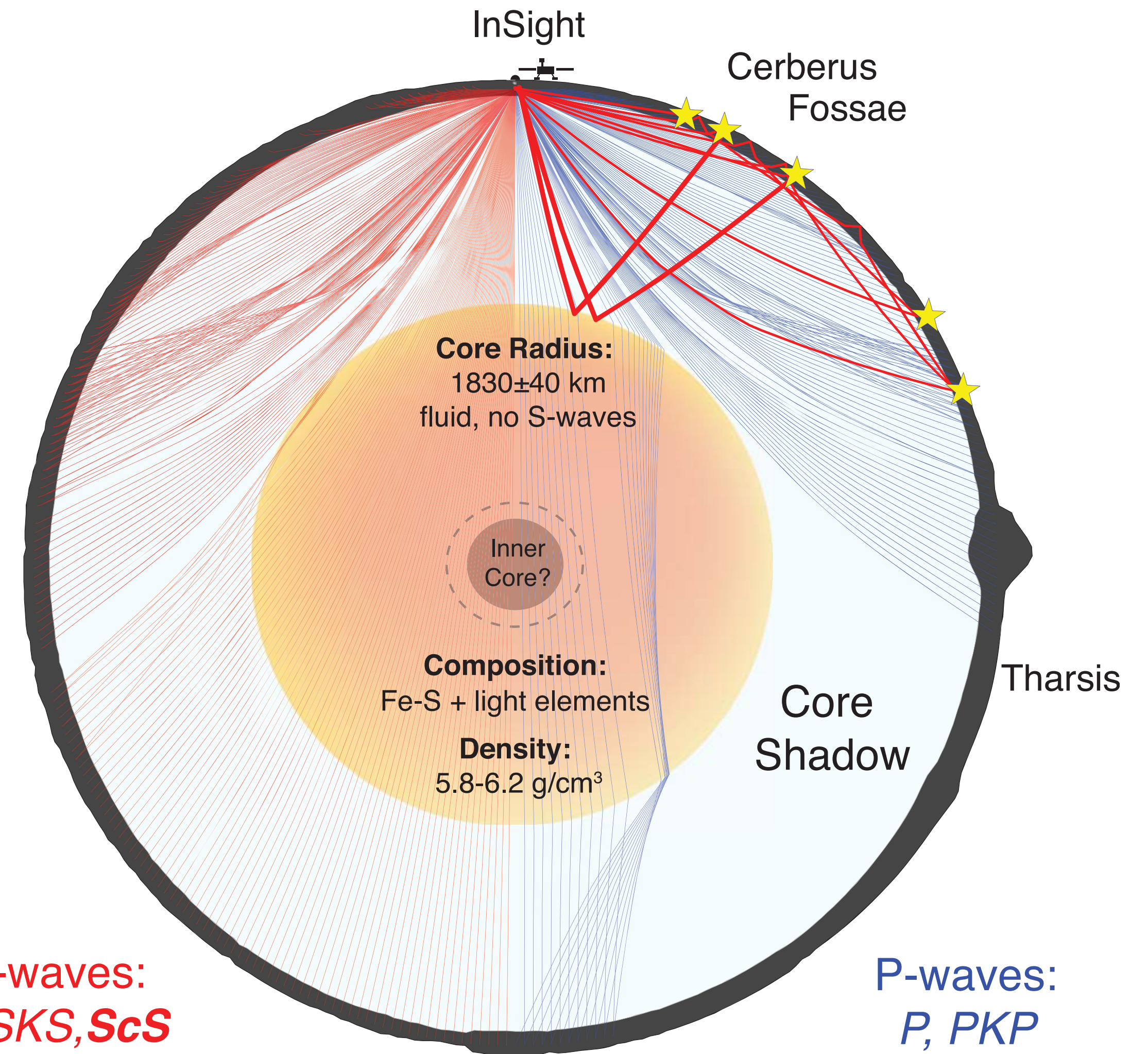
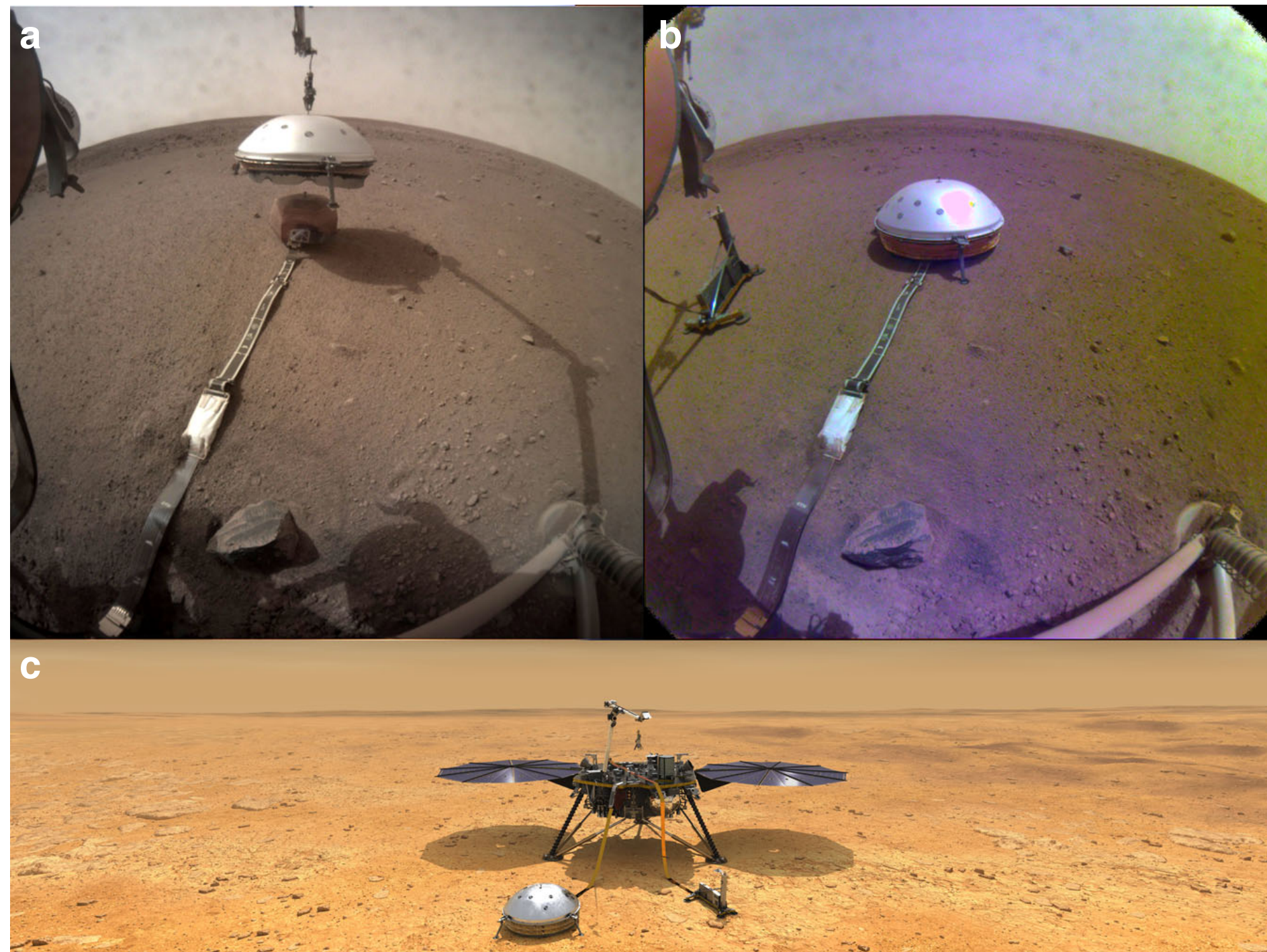


Lunar seismograph network  
Lognonné & Johnson (2015)  
in *Treatise on Geophysics 2nd Edition*

Khan & Mosegaard (2002) *J. Geophys. Res.*

# Mars' seismology

NASA's InSight and its seismograph (Image credit: NASA/JPL-Caltech)



First detection of reflection at the CMB

→ Core is large ( $R_c/R = 0.54 \pm 0.01$ ),  
and has low density ( $\rho = 5,800-6,200 \text{ kg/m}^3$ )

S-waves:  
*S, SKS, ScS*

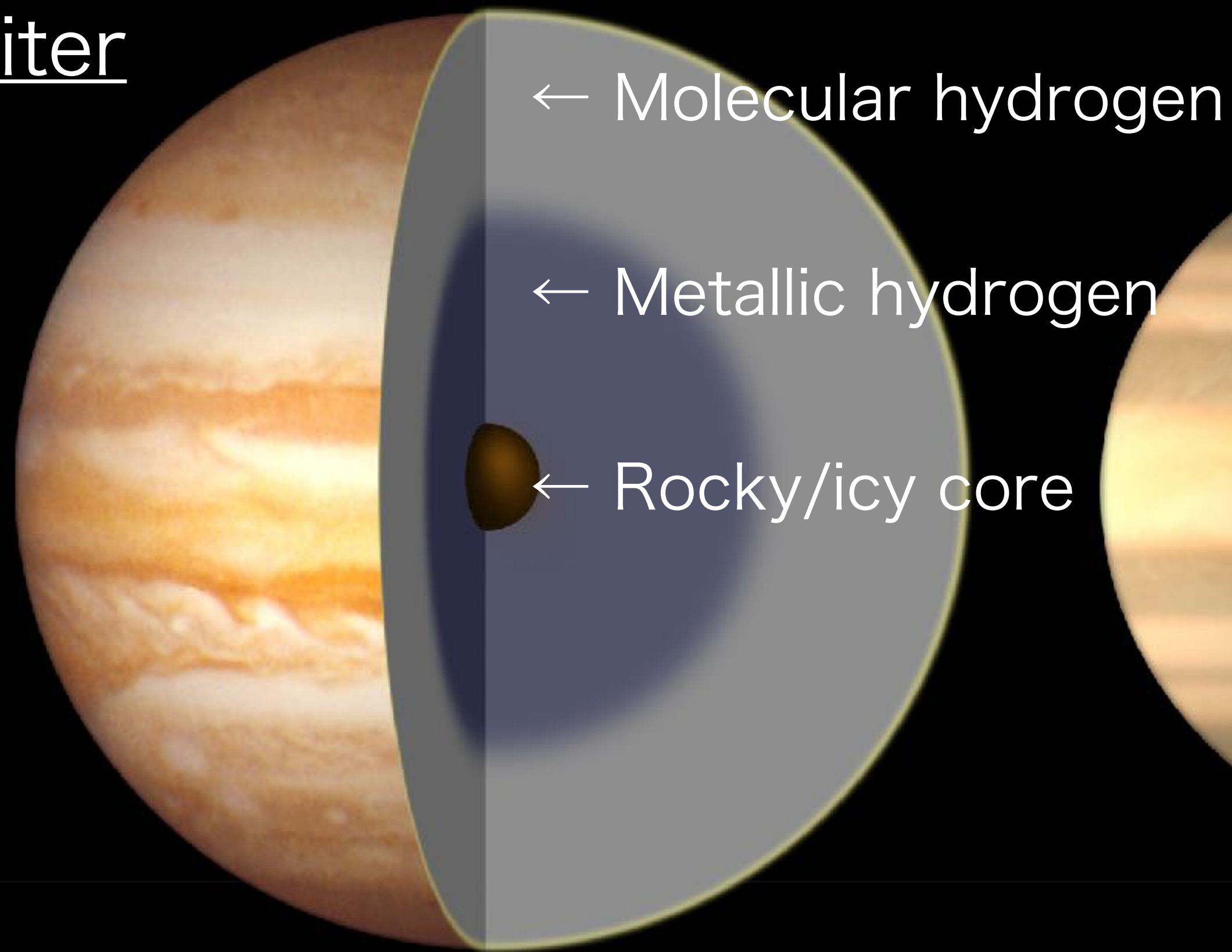
P-waves:  
*P, PKP*

Stähler et al. (2021) *Science*

# Interiors of gas giants

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Jupiter



Saturn

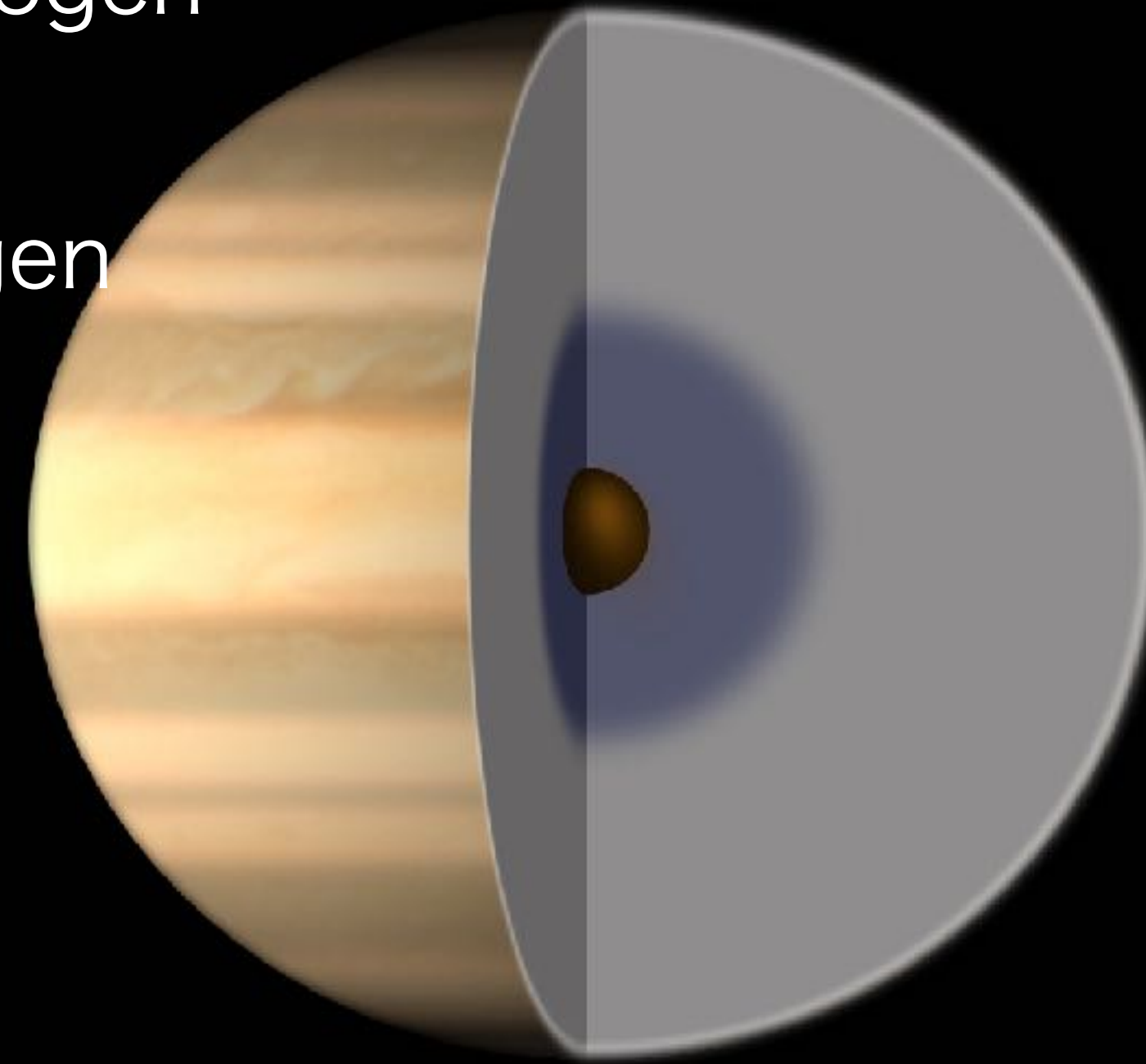
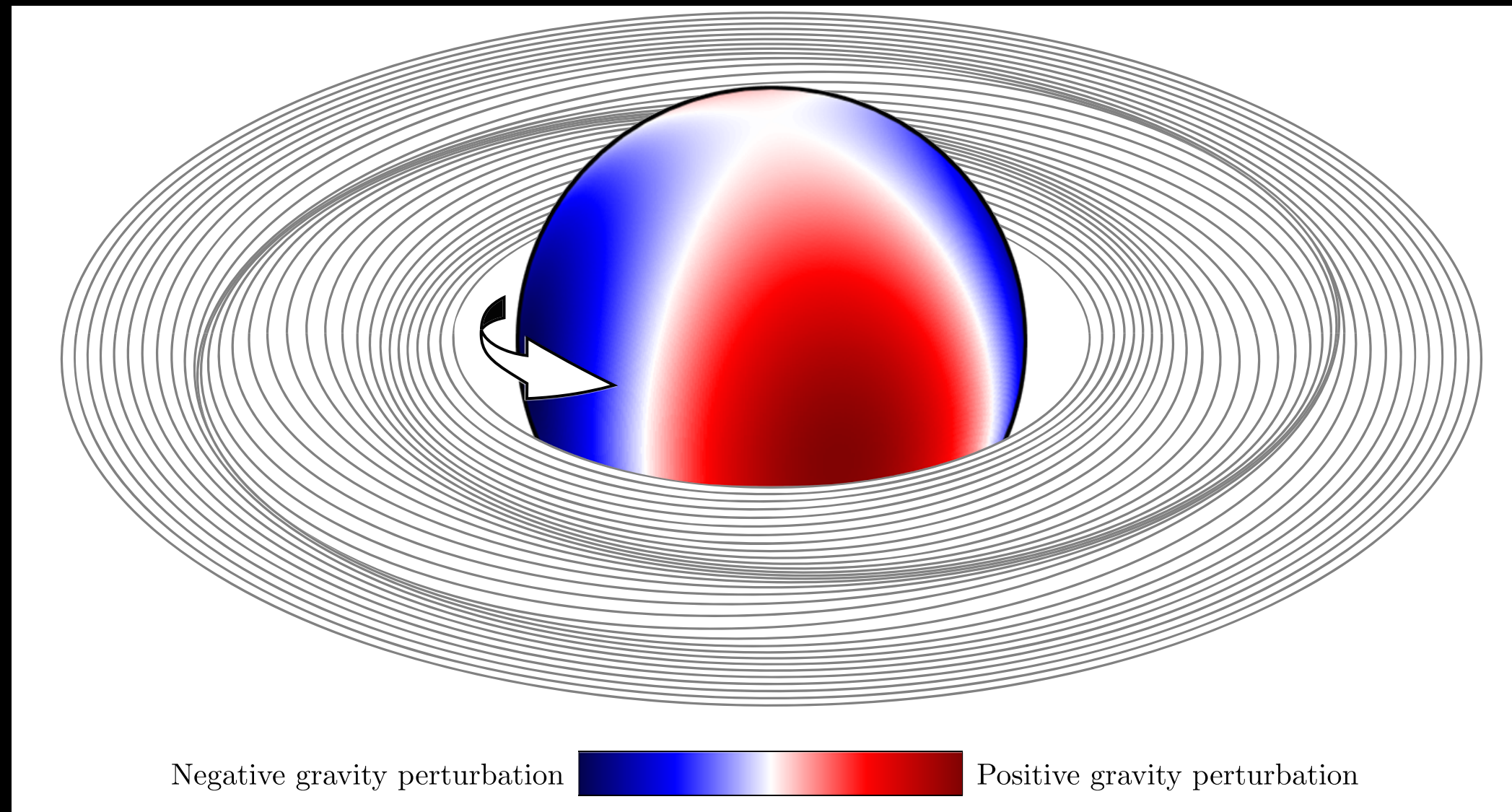


Image credit: NASA/LPI

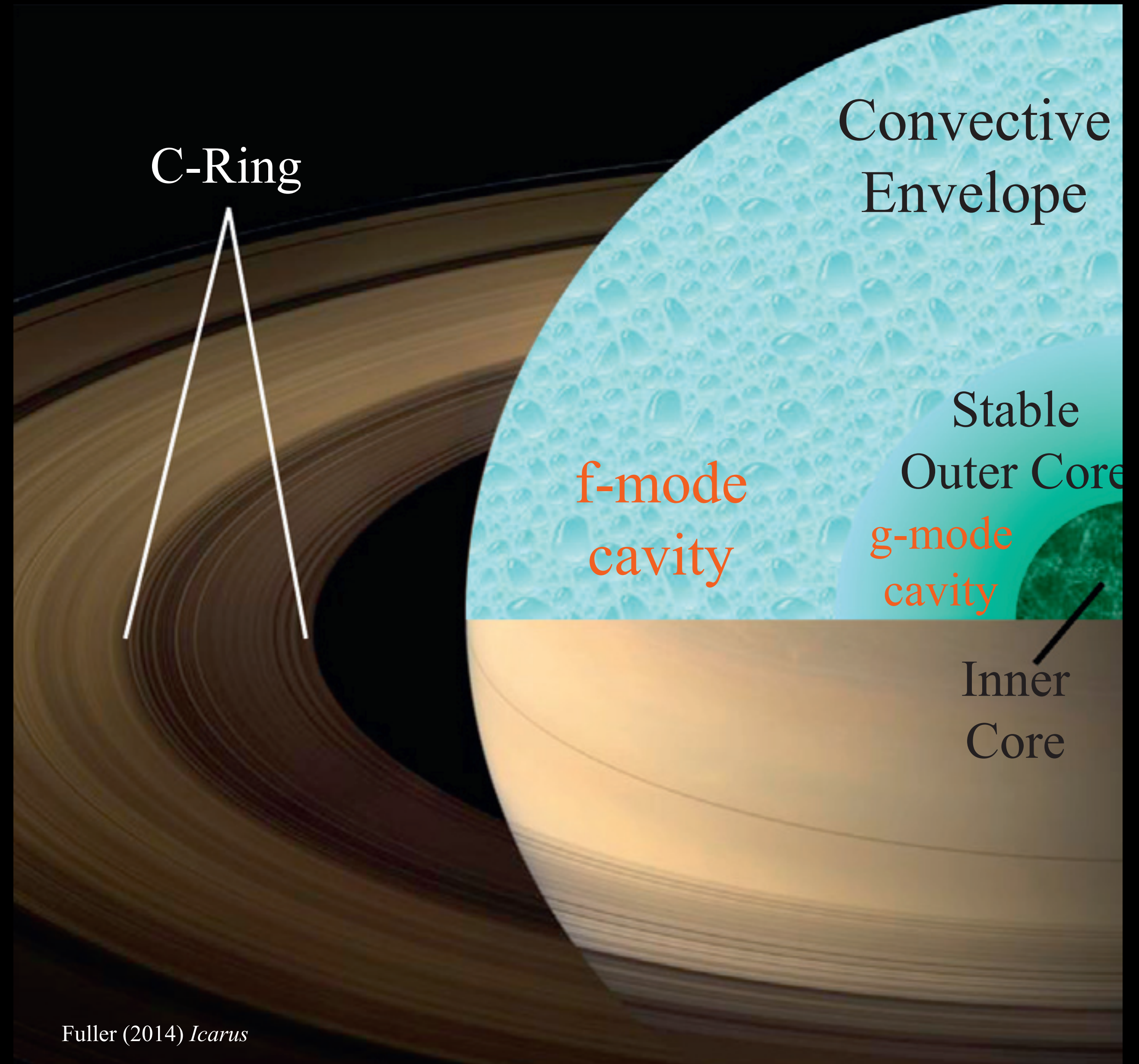
- The majority of the mass ( $\approx 70 - 90\%$ ) is hydrogen (H) and helium
- High pressure ( $>100$  GPa) interiors  $\rightarrow$  Metallic H (with free electrons)  $\rightarrow$  Magnetic dynamo

# Seismology of gas giants

Saturn's interior informed by waves on the ring



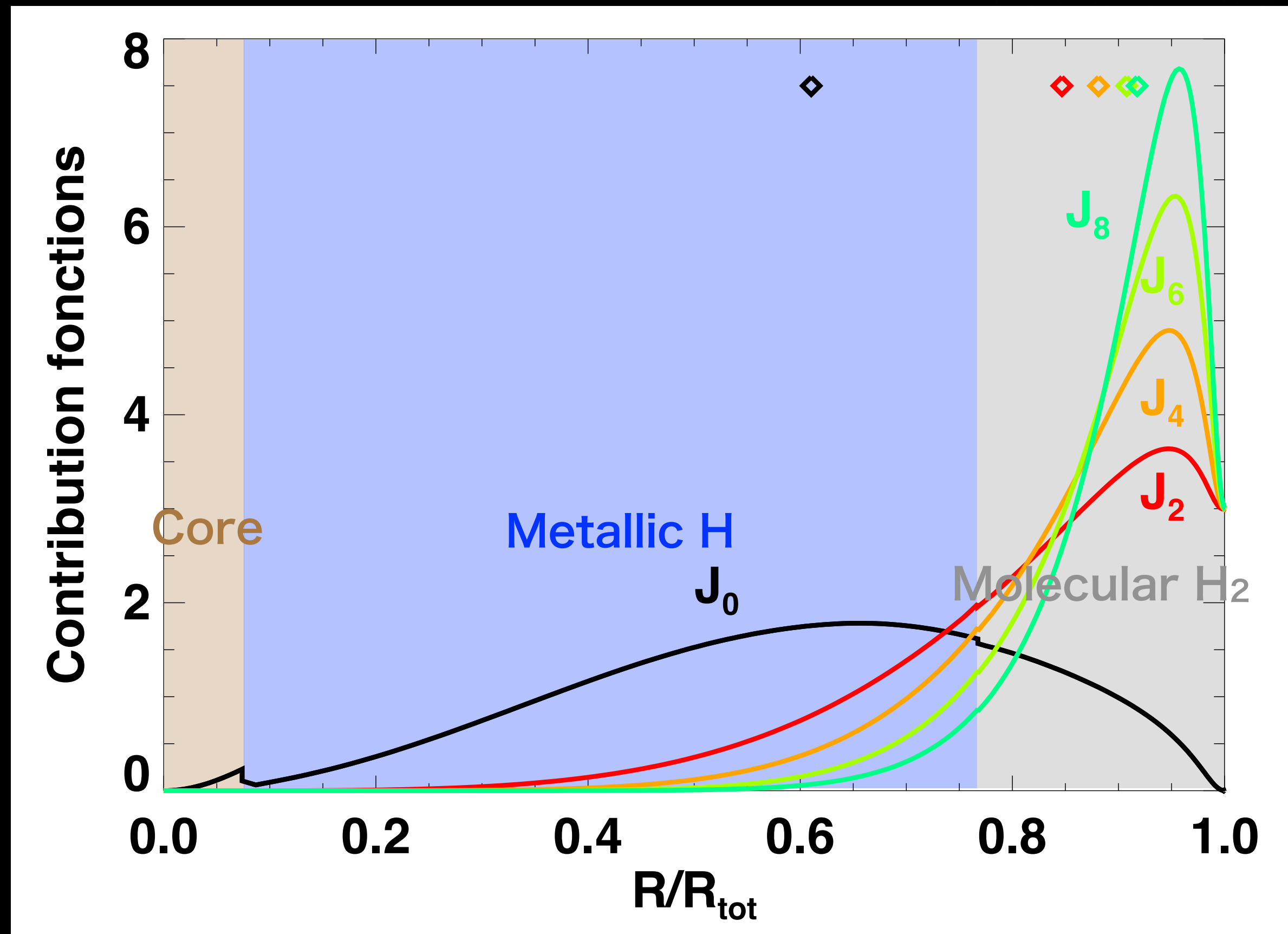
Mankovich (2020)



Fuller (2014) *Icarus*

# Gravity measurements

## Contributions to the gravitational moments of Jupiter



Guillot & Gauter (2014)

## Gravity potential

Spherical harmonic expansion

$$V = \frac{GM}{r} \left[ 1 - \sum_{n=1}^{\infty} \left( \frac{a}{r} \right)^{2n} J_{2n} P_{2n} \cos \theta \right]$$

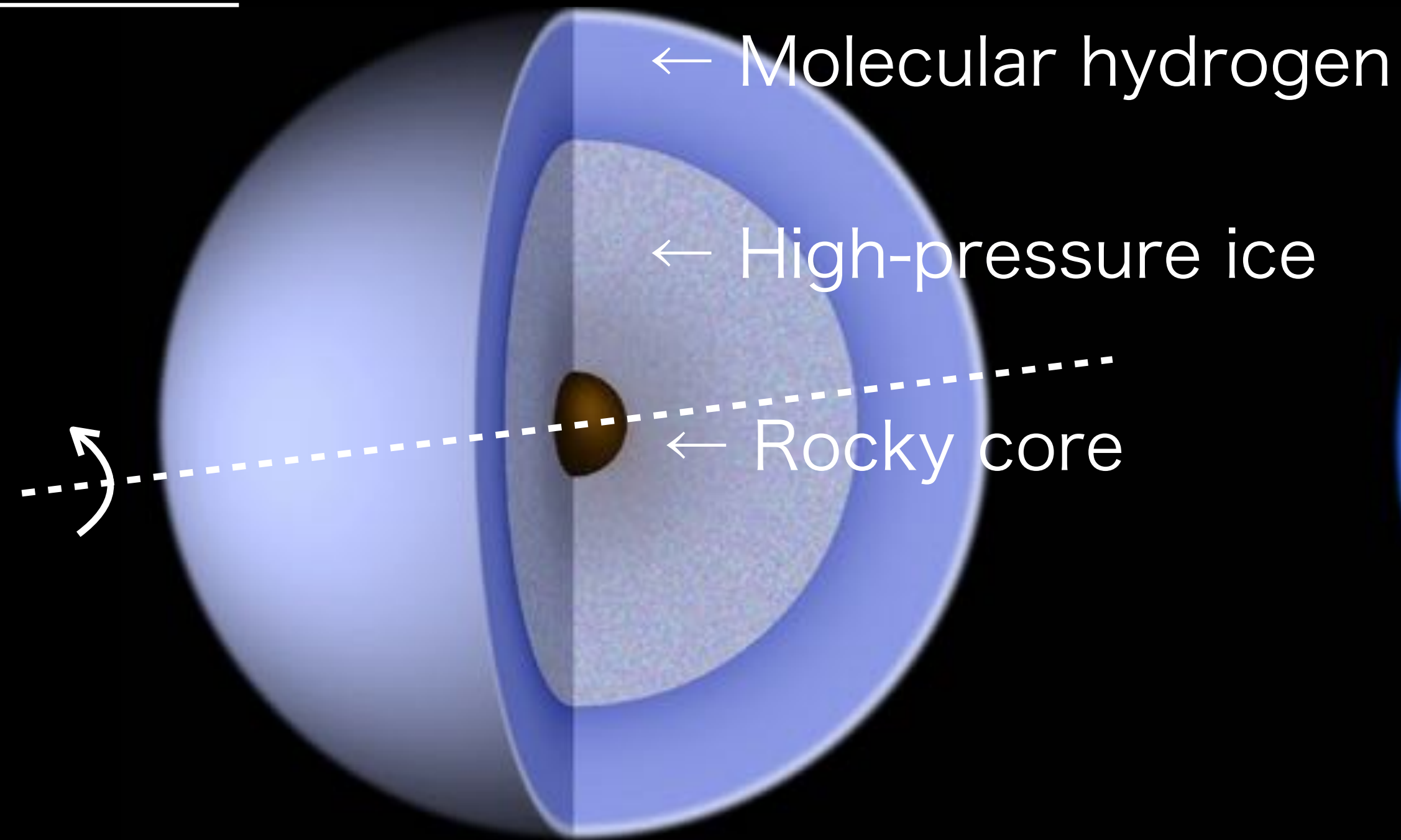




# Interiors of icy giants

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Uranus



Neptune

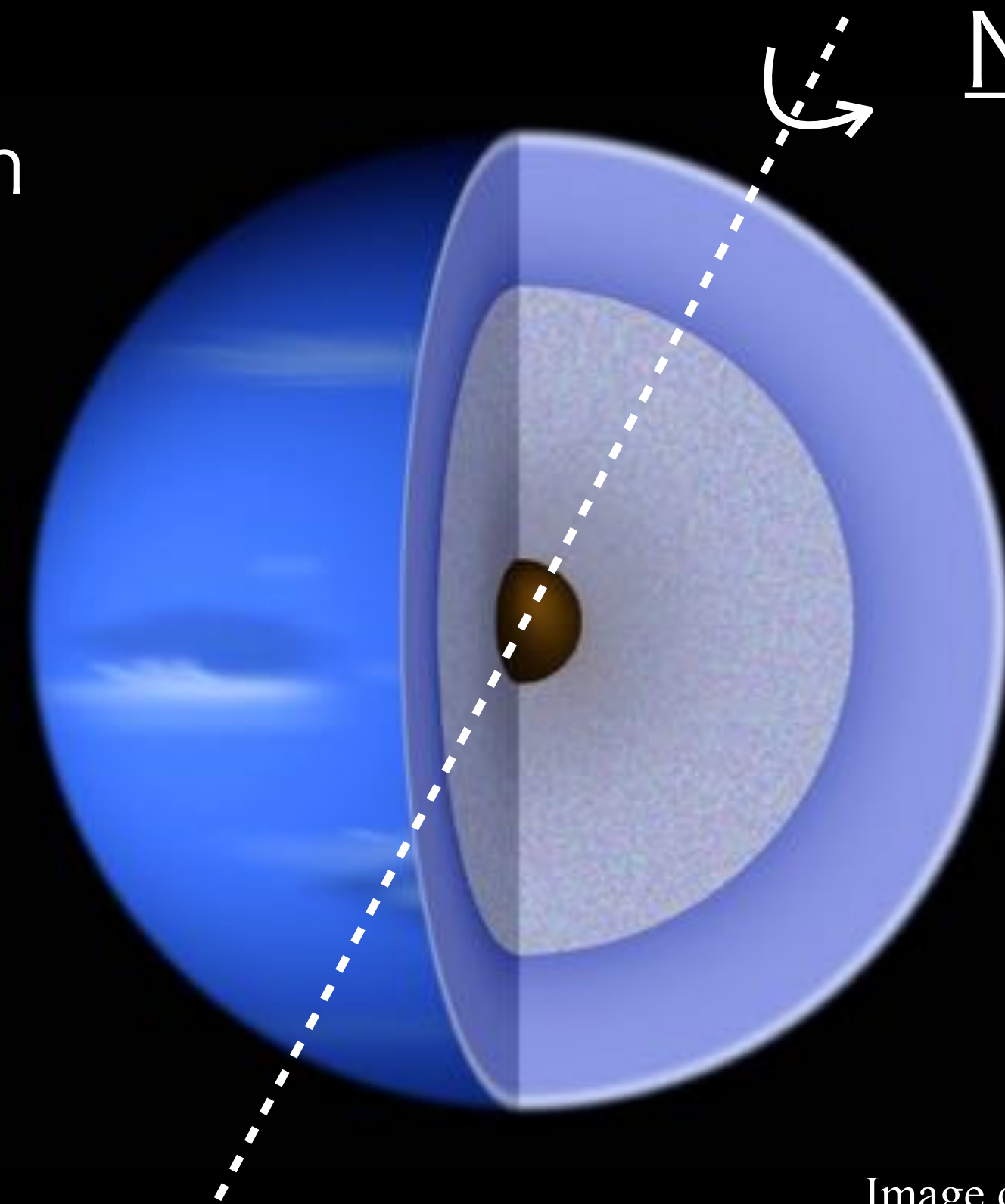


Image credit: NASA/LPI

- High-pressure ice: Super-ionic phase (protons behave like free electrons)
- Interiors less understood (no orbiter measurements so far)

# Moons

**Earth**



Radius : 1737 km

**Mars**

**Phobos**

**Deimos**

Radius : 11 km, 6.2 km

**Jupiter**



**Io**



**Europa**



**Ganymede**



**Callisto**

**Saturn**

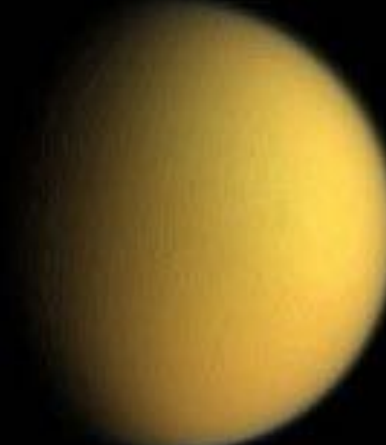
**Mimas**

**Enceladus**

**Tethys**

**Dione**

**Rhea**



**Titan**

**Hyperion**

**Iapetus**

**Phoebe**

**Uranus**

**Puck**

**Miranda**

**Ariel**

**Umbriel**

**Titania**

**Oberon**

**Neptune**

**Proteus**



**Triton**

**Nereid**

**Pluto**



**Charon**



**Earth**

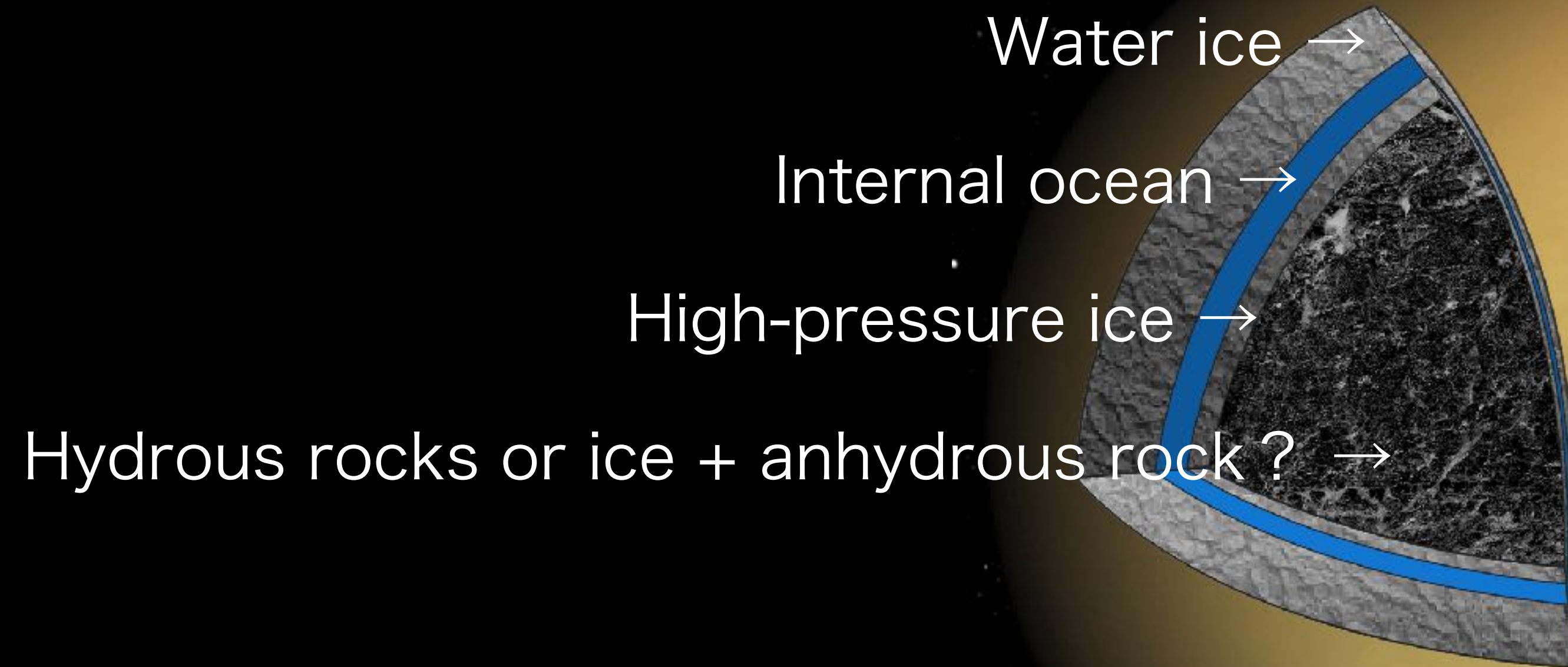
Planet	Number of moons
Mercury	0
Venus	0
Earth	1
Mars	2
Jupiter	79
Saturn	82
Uranus	27
Neptune	24

As of 2020

Image credit: NASA

# Saturn's icy moon: Titan

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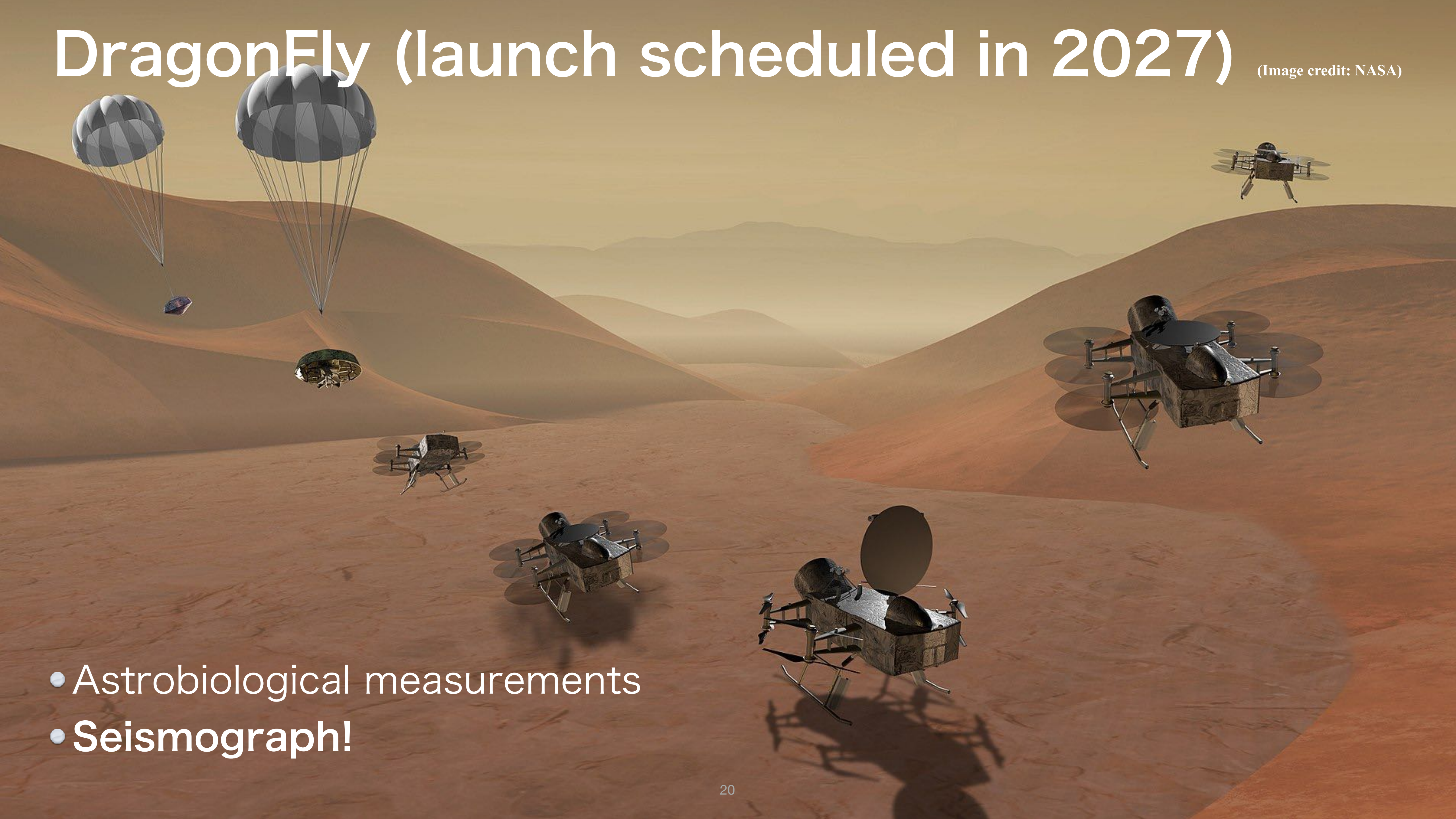


- Atmosphere:  $\sim 1.5 \times 10^5$  Pa,  $N_2$  + a few %  $CH_4$
- Photochemical haze (organic molecules)
- Lakes:  $CH_4$ ,  $C_2H_6$  liquid
- Internal ocean:  $H_2O$  (common in many icy moons)

Image credit: NASA

# DragonFly (launch scheduled in 2027)

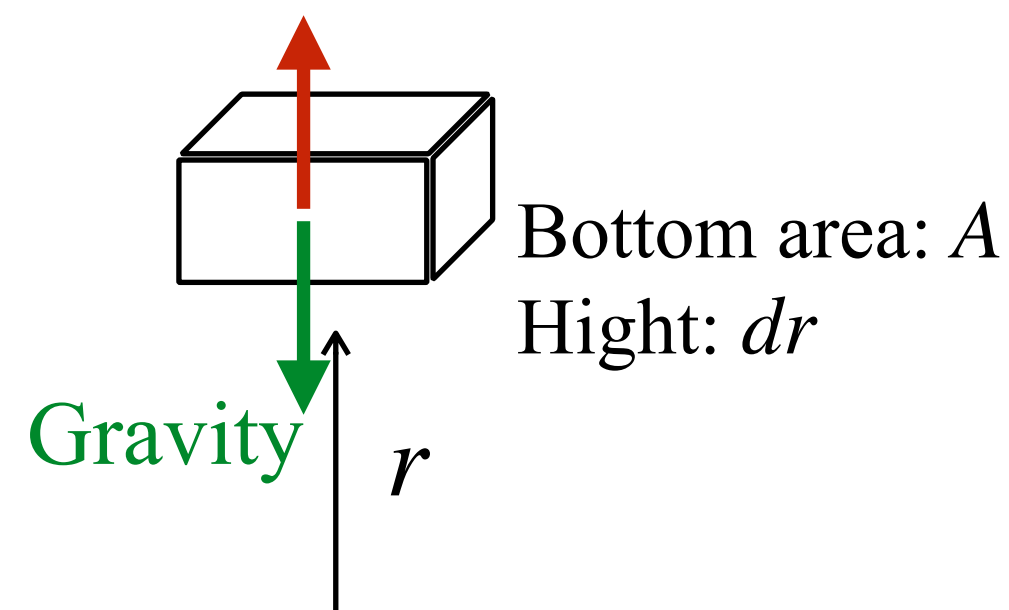
(Image credit: NASA)



- Astrobiological measurements
- **Seismograph!**

# Equation for hydrostatic equilibrium

Pressure gradient force



We consider the equilibrium of the forces.

$$\text{Pressure gradient force: } p(r)A - p(r + dr)A = -\frac{dp}{dr}dr \cdot A \quad (1)$$

$$\text{Gravity: } -\rho A dr \cdot g(r) = -\rho A dr \cdot \frac{GM_r}{r^2} \quad (2)$$

From (1) + (2) = 0, we obtain,

$$\frac{dp}{dr} = -\rho \frac{GM_r}{r^2} \quad (3) : \text{The hydrostatic equilibrium equation}$$

$p$ : pressure,  $\rho$ : density,  $g$ : gravity,

$M_r(r)$ : enclosed mass within the sphere of the radius  $r$

# An example: pressure change in the ocean

---

Let us think about pressure change when diving in the ocean.

The hydrostatic equation is given as,

$$\frac{dp}{dr} = -\rho \frac{GM_r}{r^2} \simeq \rho g \quad \text{--- (4)},$$

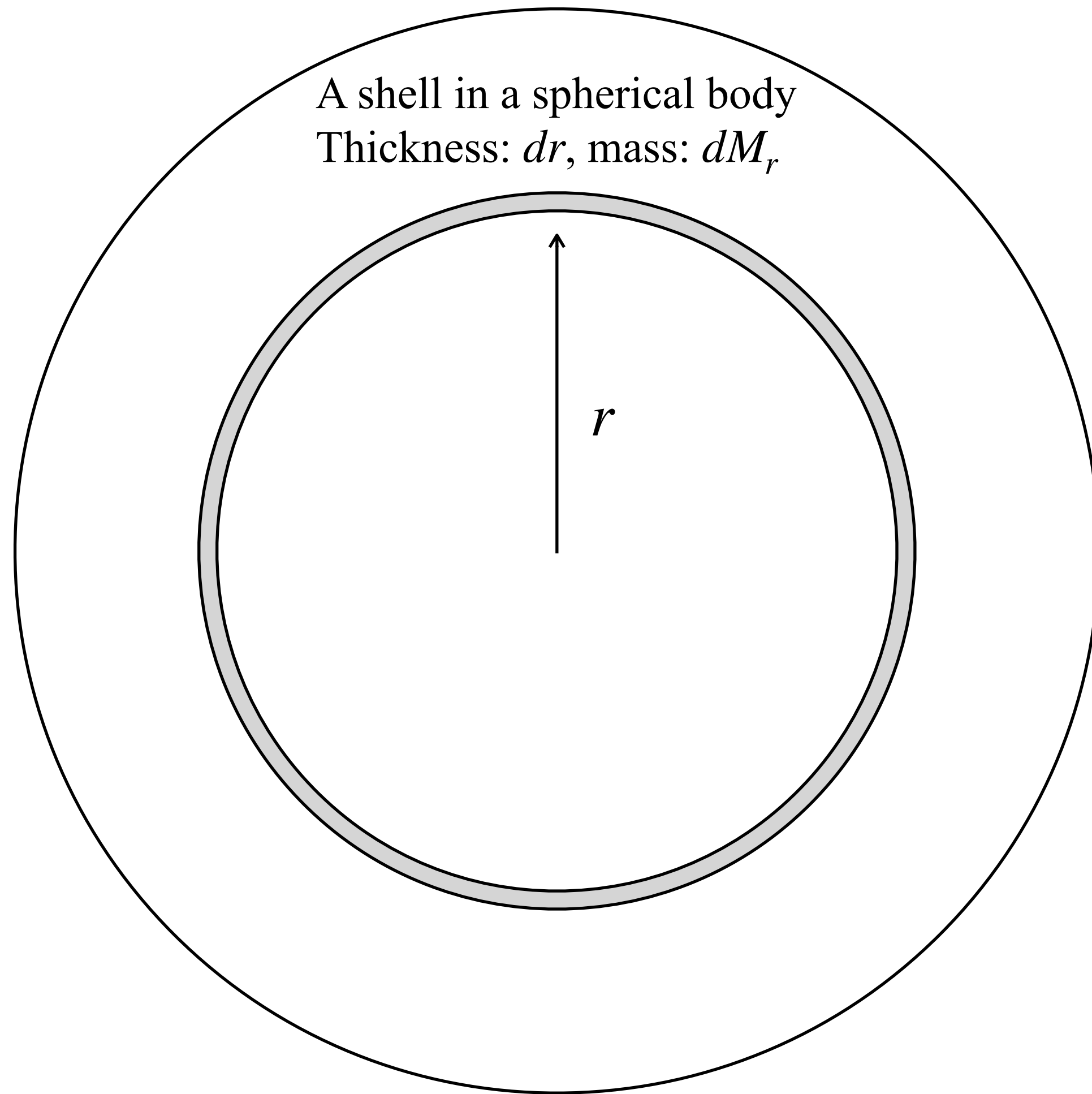
where  $g$  is the gravitational acceleration at the surface ( $9.8 \text{ m s}^{-2}$ ).

Given the density of water at 1 atm,  $\rho = 10^3 \text{ kg m}^{-3}$ , we obtain,

$$\frac{dp}{dr} = -10^3 \text{ kg m}^{-3} \cdot 9.8 \text{ m s}^{-2} \simeq -10^4 \text{ Pa m}^{-1} \simeq -1 \text{ atm/10 m} \quad \text{--- (5)}$$

# Equation for mass conservation

---



A shell in a spherical body  
Thickness:  $dr$ , mass:  $dM_r$

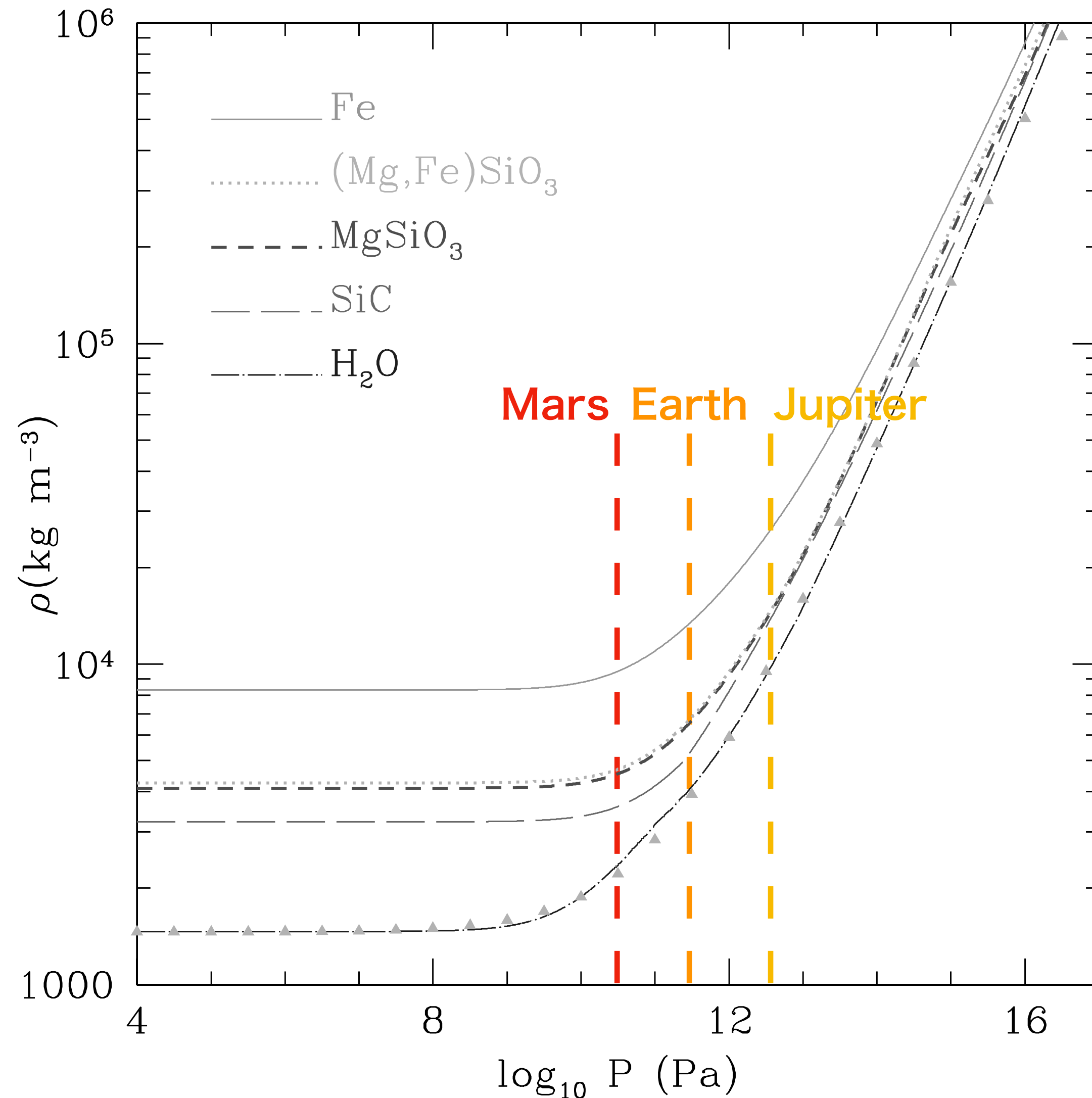
The mass of the shell,  $dM_r$ , is given by,

$$dM_r = 4\pi r^2 dr \cdot \rho \quad \text{--- (1)}$$

$$\therefore \frac{dM_r}{dr} = 4\pi r^2 \rho \quad \text{--- (2) : The mass conservation equation}$$

# Equation of state

$P - \rho$  relations for solid materials given by the equation of state



An equation of state (EoS) is a function which relates thermodynamic variables:  $p, \rho, T$

## Examples

- Ideal gas law:  $p = \frac{\rho k_B T}{m}$  — (1)

- (Third-order) Birch-Murnaghan EoS:

$$p = \frac{3}{2} K_0 (\eta^{7/3} - \eta^{5/3}) \left[ 1 + \frac{3}{4} (K'_0 - 4) (\eta^{2/3} - 1) \right] \quad \text{--- (2)}$$

Here  $\eta = \rho/\rho_0$ ,  $K$  is the bulk modulus,  $K'$  is the pressure derivative, and the subscript 0 stands for the value at the ambient conditions

Seager et al. (2007) *Astrophys. J.*



# Structure equations for a spherically-symmetric body

Equations solved for four variables –  $M_r(r)$ ,  $p(r)$ ,  $\rho(r)$ ,  $T(r)$

- Hydrostatic equilibrium:  $\frac{dp}{dr} = -\rho \frac{GM_r}{r^2}$  — (1)

- Mass conservation:  $\frac{dM_r}{dr} = 4\pi r^2 \rho$  — (2)

- Equation of state:  $p = f(\rho, T)$  — (3)

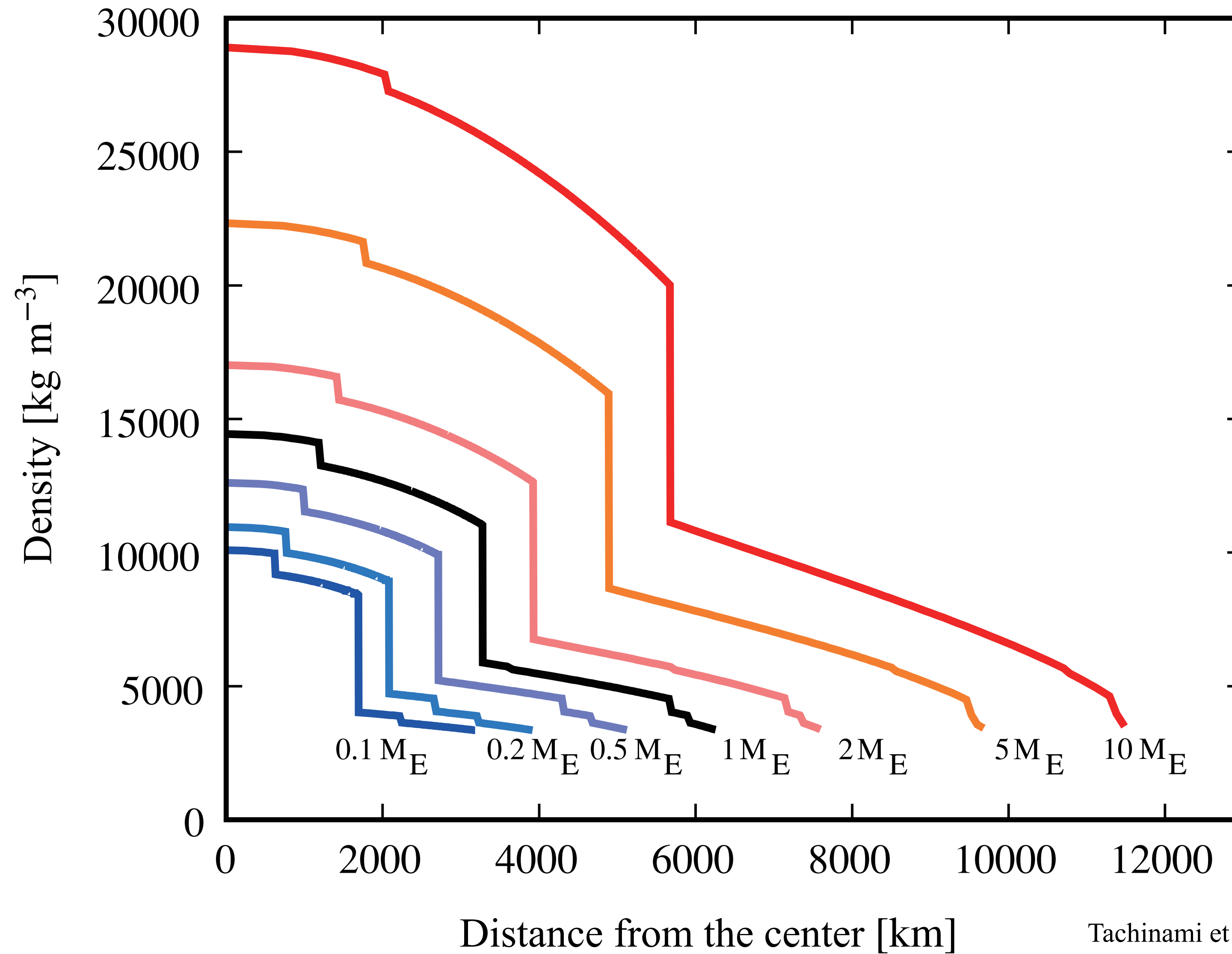
- Energy transfer:  $\frac{dT}{dr} = -\min\left(\left|\left(\frac{dT}{dr}\right)\right|_{\text{cond}}, \left|\left(\frac{dT}{dr}\right)\right|_{\text{rad}}, \left|\left(\frac{dT}{dr}\right)\right|_{\text{conv}}\right)$  — (4)

Because the temperature effect on Eq. 3 is minor for solid bodies ( $p \simeq f(\rho)$ ), Eqs. are closed with 1–3.

∴ Thermal expansion coefficient for mantle material  $\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_p \sim 10^{-5} \text{ K}^{-1}$

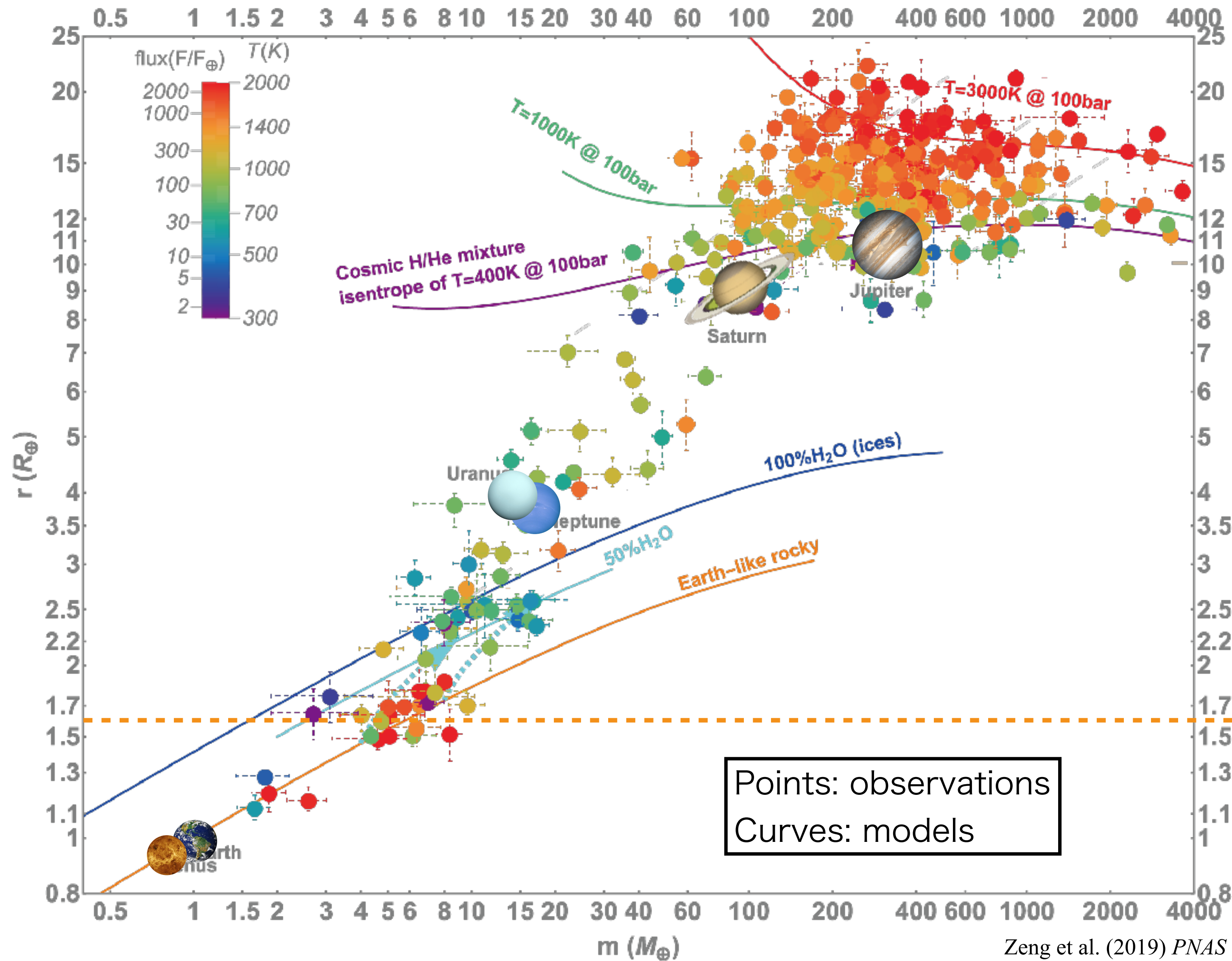
→ volume change is only  $\sim 1\% / 10^3 \text{ K}$

# The model results for Earth-like planets



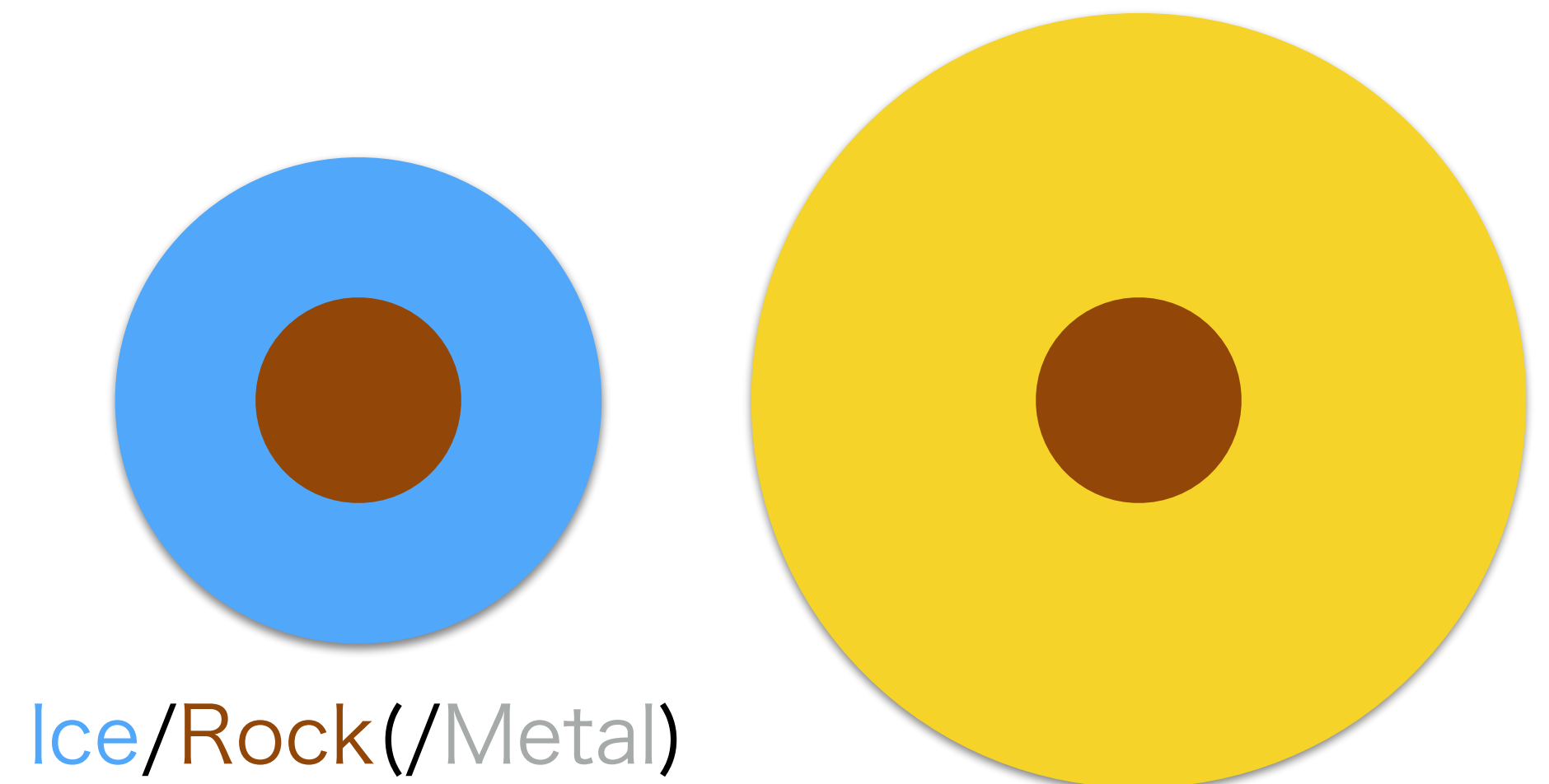
Tachinami et al. (2011) *Astrophys J.*

# Exoplanet mass-radius relations



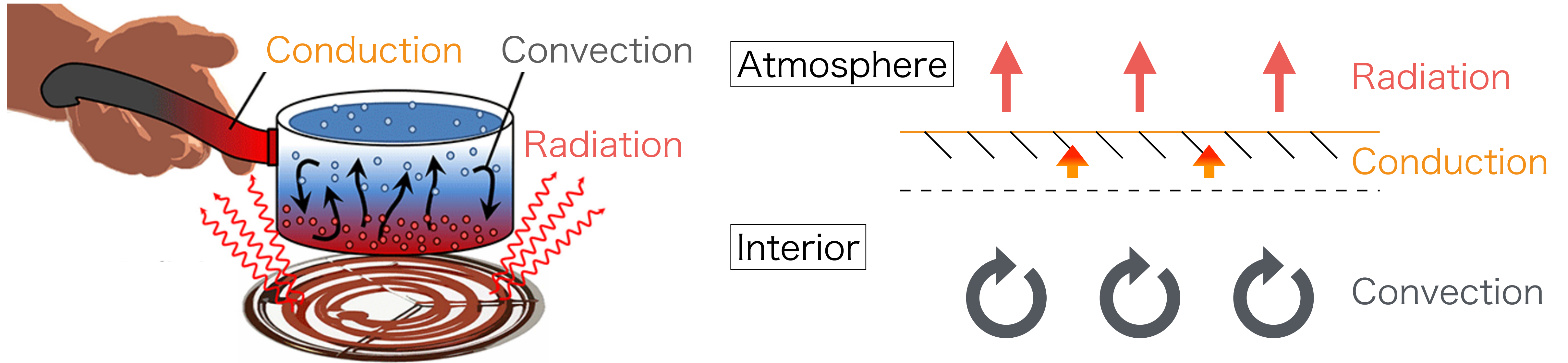
$r > 1.6 R_{\oplus}$ : ice and/or gas?

H, He/(/Ice)/Rock(/Metal)



$r < 1.6 R_{\oplus}$ : rock?

# Energy transfer in the interior



- Energy transfer in the planetary interior: convection and conduction
- Because convection cannot operate at the material interface, a conductive layer develops

# Conduction equation

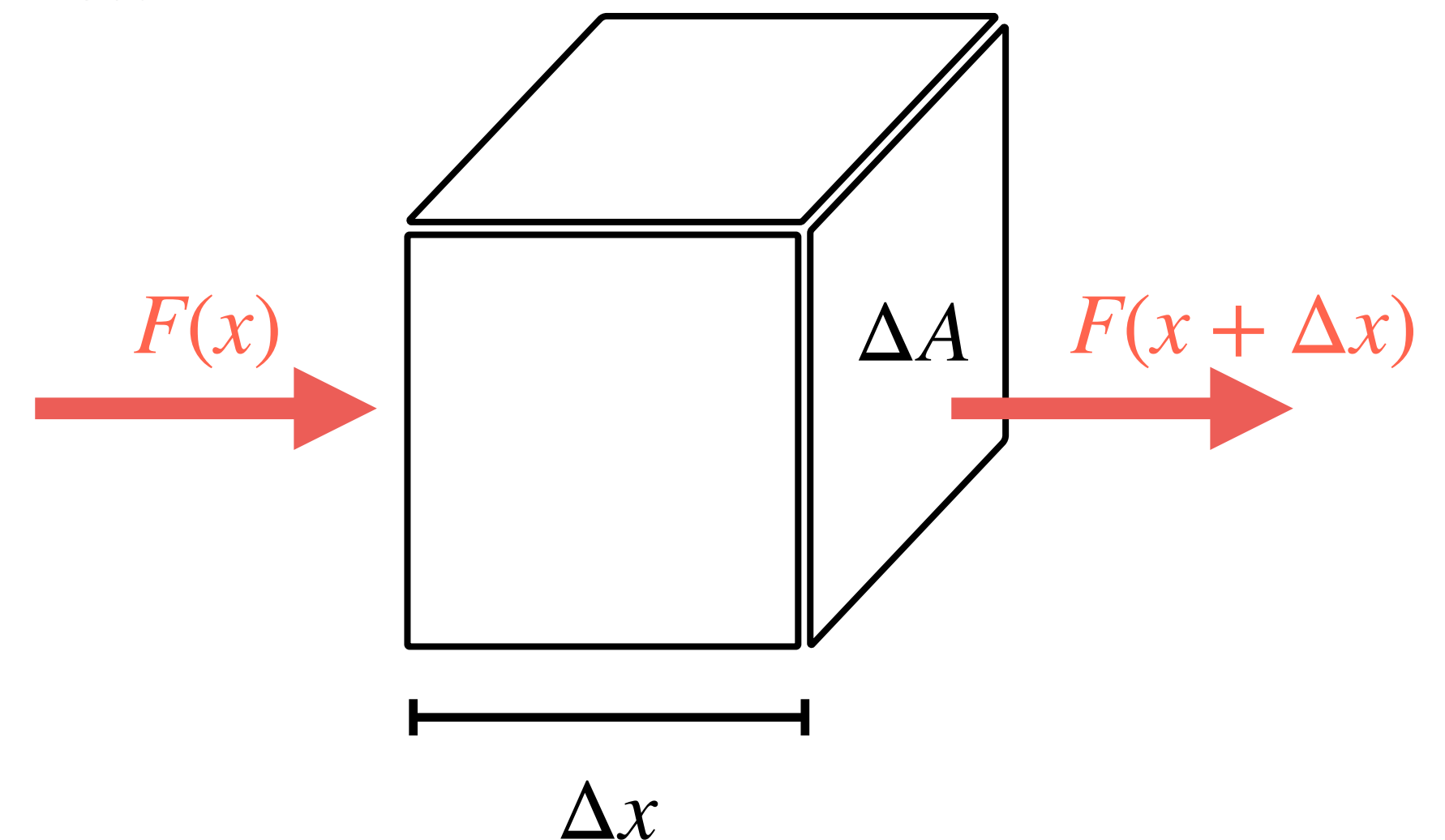
Conduction flux:  $F = -k_{\text{cond}} \frac{\partial T}{\partial x}$  — (1)

Conduction equation:  $\rho c_p \frac{\partial T}{\partial t} = -\frac{\partial F}{\partial x}$  — (2)

where  $k_{\text{cond}}$ : thermal conductivity,  $c_p$ : heat capacity (per unit mass)

When  $k_{\text{cond}}$  is constant, Eq. 2 can be written as  $\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial x^2}$  — (3)

where  $\kappa \equiv \frac{k_{\text{cond}}}{\rho c_p}$  is the thermal diffusion coefficient



# Timescale for thermal conduction

---

$$\frac{\partial T}{\partial t} = \kappa \frac{\partial^2 T}{\partial x^2} \quad \text{--- (3)}$$

Thermal diffusivity of mantle rock  $\kappa \sim 10^{-6} \text{ m}^2 \text{ s}^{-1}$  --- (4)

The distance  $l$  and its conduction timescale  $\tau$  has a relation,  $\tau \sim \frac{l^2}{\kappa}$  --- (5)

For the entire mantle  $l \sim 3 \times 10^6 \text{ m} \rightarrow \tau \sim 10^{11} \text{ year}$  !

$\therefore$  Conductive cooling is inefficient for a planet

# Rayleigh number and convective instability

- Whether a system starts to convect is determined by the Rayleigh number  $Ra$ ,

- $$Ra = \frac{\alpha \rho g \Delta T d^3}{\kappa \eta} = \text{buoyancy} / (\text{heat conduction} \cdot \text{viscosity}) \quad \text{— (1)}$$

$\alpha$ : thermal expansion coefficient,  $\rho$ : density,  $g$ : gravitational acceleration,

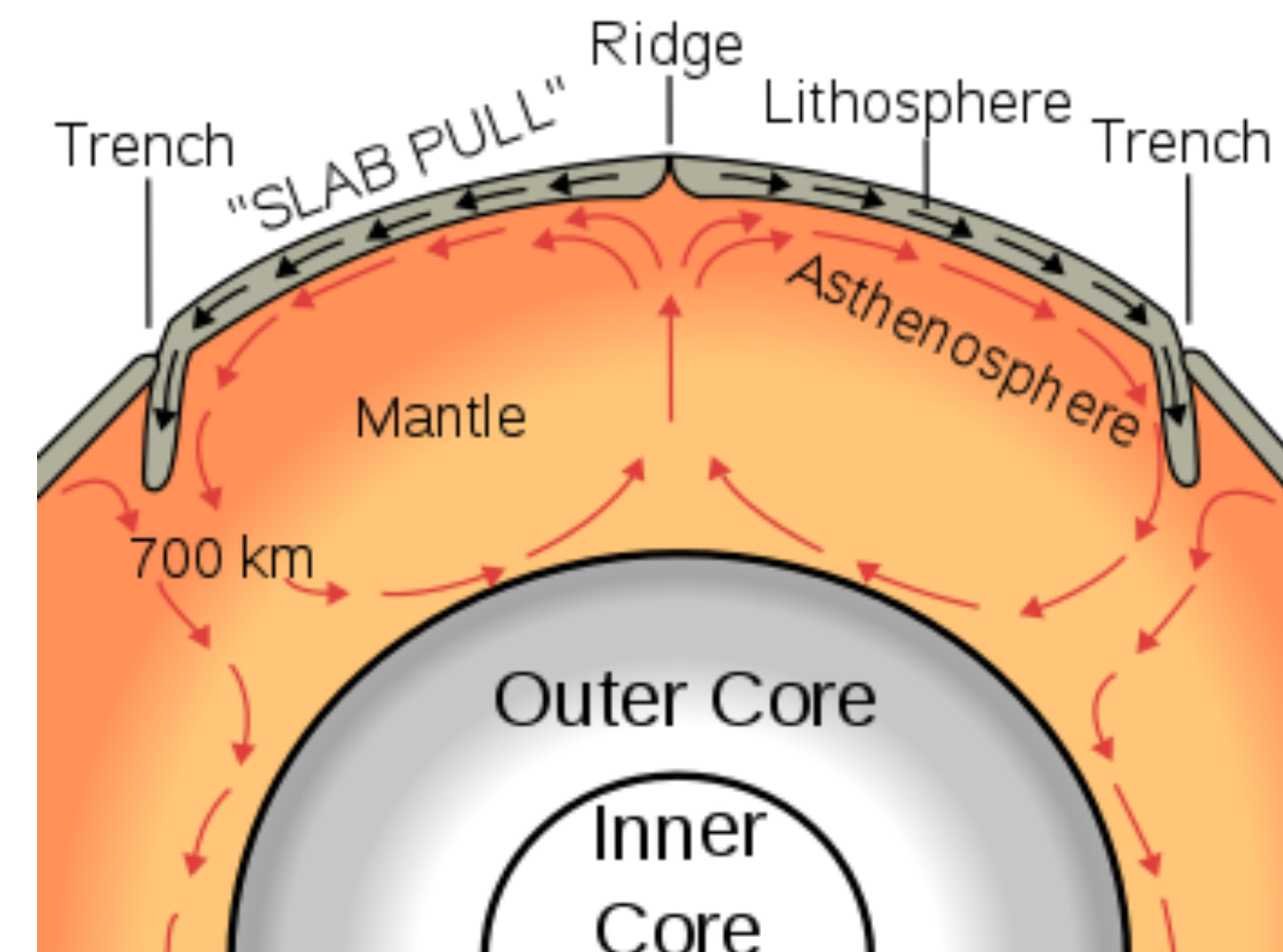
$\Delta T$ : temperature difference between the top and the bottom,

$d$ : distance from the top to the bottom,  $\kappa$ : thermal diffusion coefficient,

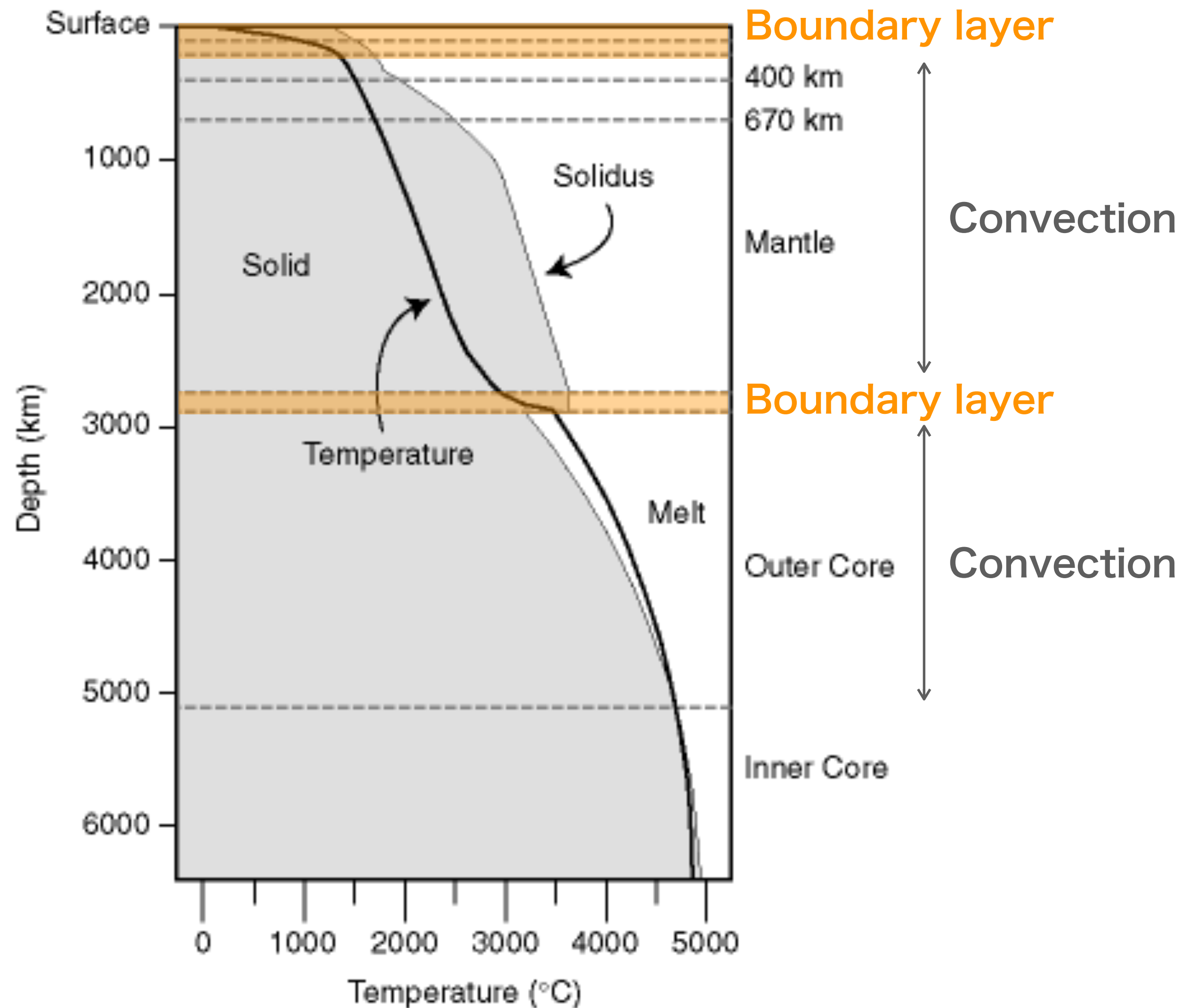
$\eta$ : viscosity coefficient

- Criterion for instability:  $Ra \gtrsim 10^3$

$\leftrightarrow$  Earth's mantle  $Ra \sim 10^7 - 10^8$



# Temperature profile of Earth's interior



- Boundary layer profile:  
Temperature gradient to conduct energy

$$\frac{dT}{dr} \simeq \left( \frac{dT}{dr} \right)_{\text{cond}} \equiv -\frac{F_{\text{int}}}{k_{\text{cond}}} \quad (1)$$

Thermal conductivity:  $k_{\text{cond}} \equiv \rho C_p \kappa$

- Convective layer profile:  
Adiabatic lapse rate

$$\frac{dT}{dr} \simeq \left( \frac{dT}{dr} \right)_{\text{ad}} = -\frac{\alpha g T}{C_p} \quad (2)$$



# Derivation of adiabatic lapse rate (for physics students)

---

$$\text{Entropy change: } dS = \left(\frac{\partial S}{\partial T}\right)_p dT + \left(\frac{\partial S}{\partial p}\right)_T dp = \frac{C_p}{T} dT + \left(\frac{\partial S}{\partial p}\right)_T dp \quad \text{--- (1)}$$

$$\text{Gibbs free energy change } dG = d(U + pV - TS) = Vdp - SdT \quad \text{--- (2),}$$

$$S = - \left(\frac{\partial G}{\partial T}\right)_p, \quad V = \left(\frac{\partial G}{\partial p}\right)_T \quad \text{--- (3).}$$

$$\therefore \left(\frac{\partial S}{\partial p}\right)_T = - \left(\frac{\partial}{\partial p} \left(\frac{\partial G}{\partial T}\right)_p\right)_T = - \left(\frac{\partial}{\partial T} \left(\frac{\partial G}{\partial p}\right)_T\right)_p = - \left(\frac{\partial V}{\partial T}\right)_p \quad \text{--- (4). (Maxwell relations)}$$

$$\text{Substituting Eq. 4 into Eq. 1, we obtain } dS = \frac{C_p}{T} dT - \left(\frac{\partial V}{\partial T}\right)_p dp \quad \text{--- (5)}$$

$$\text{Finally, substituting } \alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_p \quad \text{--- (6) into Eq. 5, we obtain } dS = \frac{C_p}{T} dT - \alpha V dp. \quad \text{--- (7)}$$

$$\therefore \left(\frac{\partial T}{\partial p}\right)_S = \frac{\alpha T}{\rho C_p} \quad \text{--- (8)} \rightarrow \left(\frac{dT}{dz}\right)_{\text{ad}} = \frac{dp}{dz} \left(\frac{\partial T}{\partial p}\right)_S = - \frac{\alpha g T}{C_p} \quad \text{--- (9)}$$

# Summary

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- How to know planetary interior structures:
  - bulk density, seismology, gravity measurements, etc.
- Earth's interior: crust, mantle, outer and inner core
- Diversity in planetary compositions: rocky, gaseous, icy
- Planetary structure equations:
  - Hydrostatic equilibrium
  - Mass conservation
  - Equation of state
  - Energy transfer: conduction, convection, radiation

# Report assignment

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Summarize your answers into a short report and submit it by the beginning of the next lecture (either directly, to my post-box, or by e-mail to [hiro.kurokawa@elsi.jp](mailto:hiro.kurokawa@elsi.jp)).

1. Thermal conduction determines the temperature profile in the boundary layer. Using the physical quantities given below, estimate the temperature gradient in the upper boundary layer of Earth's interior (the top  $\sim 100$  km). Answer with one significant digit.

$$F_{\text{int}} = 0.09 \text{ W} \cdot \text{m}^{-2}, \kappa \simeq 1 \times 10^{-6} \text{ m}^2 \cdot \text{s}^{-1}, \rho \simeq 3 \times 10^3 \text{ kg} \cdot \text{m}^{-3}, c_p \simeq 1 \times 10^3 \text{ J} \cdot \text{kg}^{-1} \cdot \text{K}^{-1}$$

$$\therefore \left( \frac{dT}{dz} \right)_{\text{cond}} = - \frac{F_{\text{int}}}{k_{\text{cond}}} = - \frac{F_{\text{int}}}{\rho C_p \kappa} \simeq - \boxed{\phantom{00}} \text{ K} \cdot \text{km}^{-1}$$

2. Let's assume that you are a hot-spring (*onsen*) enthusiast and want to dig for a hot spring of your own. Using the result of Q1, discuss how deep you need to dig a hole in the ground.